Angular scattering of the Sahara dust aerosol

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Abstract. Soil erosion aerosols can be transported considerable distances, the Sahara being one of the major sources in the world. In June 2016 the volume scattering function of the atmospheric aerosol was determined in the Sierra Nevada, Spain, at an altitude of 2500 m. Measurements were performed with a polar nephelometer permitting measurements between scattering angles of 5 to 175°. The values at the missing angles could be estimated to a high accuracy, using the shape of the scattering function adjacent to the missing angles, and thus a complete volume scattering function was available. During the measuring period intrusions of long-range transported Sahara aerosol happened several times. The classification of the aerosol was done by back trajectories and by the Ångström exponent of the wavelength-dependent scattering coefficient, which was determined by a three-wavelength Integrating Nephelometer. The phase function of the Sahara aerosol had a stronger forward scattering, and less backscattering compared to the non-Sahara aerosol, which is in agreement with other findings for irregular particles. The asymmetry parameter of the phase function is the best characteristic to distinguish Sahara aerosol from non-Sahara aerosol. In this study the asymmetry parameter for the Sahara aerosol was larger than 0.65, whereas the non-Sahara aerosol had an asymmetry parameter below 0.6. A comparison with measurements performed with long-range transported Gobi desert aerosols observed in Kyoto, Japan, showed very similar results.

1 Introduction

Deserts are a major source of aerosol particles. On a global scale desert aerosol contributes 60–1800 Tg yr⁻¹ of the total yearly aerosol production of 2900–4000 Tg (Jaenicke, 1988). Junge (1979) estimates the global source strength of deserts as 260 to 400 Tg yr⁻¹, with the Sahara contributing 60 to 200 Tg yr⁻¹. Zender et al. (2004), using all estimates published between 2001 and 2004, give a range of global dust emissions of between 1000 and 2150 Tg yr⁻¹. Recently the total PM₁₀ (particles smaller than 10 µm) emission of desert dust particle emissions was estimated as 1700 Tg yr⁻¹; the atmospheric loading of dust particles PM₁₀ is estimated as 20 Tg (Kok et al., 2017). For PM₂₀ (particles < 20 µm) the values are 3000 Tg yr⁻¹ and 23 Tg. The clay-sized particles (diameter smaller than 2 µm) are an important subgroup due to their efficient light scattering, but account only for 3.5 % to 5.7 % of the PM₁₀ emissions.

In contrast to most of the other particles present in the atmosphere, desert aerosol particles are produced from minerals by mechanical processes, which obviously led to irregularly shaped particles. This is documented by innumerable electron micrograph studies (see e.g. Falkovich et al., 2001; Iwasaka et al., 2003; Kandler et al., 2007, 2011). The mechanical production of particles results in particle sizes usually larger than 1 µm. Mean sizes range up to 10 µm (Kontratyev et al., 2006; Alfaro et al., 2003; Cheng et al., 2005; Xin et al., 2005; Alfaro and Gomez, 2001; Falkovich et al., 2001), and during atmospheric transport the size distribution is modified, mainly due to the shorter lifetime of larger particles: the average residence time in the troposphere is estimated as 10 days for 1 µm particles and 3 days for 10 µm particles (Jaenicke, 1988). Other estimates of the lifetime of
the size-dependent desert particles are 11 days for 1 µm and 0.4 days for 15 µm (Kok et al., 2017). Obviously the particle size distribution is modified when transported in the atmosphere, since larger particles have a considerable sedimentation velocity; therefore, the clay-sized particles amount to 15 % of the global atmospheric load, while accounting only for 4.3 % of the emitted mass. With the global wind system the smaller Sahara aerosol particles can be transported considerable distances, have been observed in the Amazonas basin (Formenti et al., 2001) and are a major source of mineral supply to the Amazonas or Congo basins (Okin et al., 2004).

Among many other effects, the desert aerosol is expected to have an influence on the radiative balance; see e.g. Obregón et al. (2015), Valenzuela et al. (2012a, b), and Antón et al. (2014). Due to the irregular shape of the particles, they have a larger surface-to-volume ratio, leading to a higher extinction efficiency compared to volume-equivalent spherical particles (Kocijáň and Horváth, 2005; Kok et al., 2017). The maximum efficiency is around an equivalent diameter of 1 µm; thus, the clay-sized particles are especially important. Recently it has been found that the cooling effect of desert aerosol might be smaller than expected (Kok et al., 2017); in the source region coarse dust particles (> 5 µm) dominate and they absorb both solar and thermal radiation; furthermore, the effect is augmented by the bright surface of the desert; distant from source regions the cooling effect of the fine dust seems to be less than previously anticipated.

The dust particles have an irregular shape; thus, their scattering properties are difficult to model. If the shape, size, and possible inhomogeneities of the particles are known, the scattering matrix of the particle could be calculated (Lindqvist et al., 2014). Since it is not possible to have an electron microscope graph of every particle, simplifications are used: Lindqvist et al. (2014) considered the following models: homogeneous particles of irregular shape, Gaussian random spheres, spheroids and spheres. Whereas the scattering properties of homogeneous irregular particles do not differ much from the reality of inhomogeneous particles, the other models give considerable deviations; e.g. when assuming spheres instead of irregular chrysotile particles, the asymmetry parameter is 0.66 instead of 0.82, and the lidar ratio is 4 instead of 48 sr.

Obviously this discrepancy cannot be tolerated. Thus several indirect methods are in use, such as inverting sun and sky radiation measurements (e.g. Estellés et al., 2007), using measured size distributions and Mie theory to calculate the scattering, or using an integrating nephelometer with two scattering ranges and a Henyey–Greenstein model phase function (Andrews et al., 2006). Model assumptions always have to be made in order to obtain the complete scattering properties of the aerosol containing irregularly shaped particles, which may create considerable uncertainties as shown above. But the exact scattering properties, especially the asymmetry parameter, are urgently needed to estimate the effect on climate. Below we report the first direct measurements of the complete volume scattering function of the Sahara aerosol transported to the Iberian Peninsula, which were performed during several dust outbreaks. No assumption or model is needed in order to obtain the volume scattering function and the asymmetry parameter.

2 Definitions, units, and nomenclature

The volume scattering function γ(θ) of the aerosol is defined as follows. See Fig. 1: let a volume dV of aerosol be illuminated by radiation with a flux density S. The light flux dΦ scattered into a solid angle dω at the scattering angle θ is obtained as

\[ d\Phi = S \cdot \gamma(\theta) \cdot d\omega \cdot dV. \]

The unit is [γ] = m⁻¹ sr⁻¹. The total scattering coefficient \( \sigma_s \) is obtained by integrating the volume scattering function \( \gamma(\theta) \) over a whole solid angle:

\[ \sigma_s = \int_0^{4\pi} \gamma(\theta) d\omega, \]

or for scattering with rotational symmetry with respect to the incident beam:

\[ \sigma_s = 2\pi \int_0^{\pi} \gamma(\theta) \sin(\theta) d\theta. \]

The scattering coefficient can be understood as the fraction of the light flux scattered per unit length out of a parallel beam of light; its unit is [σs] = m⁻¹. Both the volume scattering function and the scattering coefficient are extensive properties.

The phase function is an intensive property and describes the relative angular dependence of the scattered light of a volume element of particles. For the phase function \( P(\theta) \) we have used the following definition: \( P(\theta) = 4 \cdot \sigma_s \cdot \gamma(\theta) / \sigma_s \).

The angular distribution of the scattered light frequently is characterized by two parameters.
The asymmetry parameter \( g \) is obtained by folding the phase function with \( \cos(\theta) \); therefore,
\[
g = \frac{1}{2} \int_{0}^{\pi} P(\theta) \sin(\theta) \cos(\theta) d\theta.
\]

The asymmetry parameter is zero for symmetric scattering such as Rayleigh scatter of the air molecules and \( g = 1 \) for only forward scatter. Another characteristic is the fraction \( b \) of the backscattered radiation. It is obtained by integrating
\[
\frac{1}{2} P(\theta) \sin(\theta) \] between \( 1/2\pi \) and \( \pi \).

Both parameters \( g \) and \( b \) are intensive. The lidar ratio, \( S \), is defined as the ratio of the extinction coefficient, \( \sigma_e \), and the volume scattering function at 180°, \( \gamma(180°) \), i.e. \( S = \sigma_e / \gamma(180°) \), or \( S = 4\pi / (\gamma(180°) \cdot \omega) \), with \( \omega \) the single scattering albedo, i.e. the ratio of the scattering coefficient to the extinction coefficient. The single scattering albedo cannot be measured directly with the polar nephelometer. We have used values for \( \omega \), which were obtained by inverting data from sun and sky photometers during this study. The average for the Sahara aerosol was \( \omega = 0.928 \) and \( \omega = 0.943 \) for the non-Sahara aerosol. These values are in agreement with previous findings (Valenzuela et al., 2012a, b).

The measurement of the scattering coefficient at three wavelengths (e.g. with an integrating nephelometer, Charlson et al., 1967) permitted the determination of the wavelength dependence of the scattering coefficient. In many cases it can be represented by a power law relation
\[
\sigma(\lambda) = \sigma(\lambda_0) \cdot (\lambda/\lambda_0)\alpha
\]
with \( \alpha \) the Ångström exponent (Ångström, 1929, 1930). The value of \( \alpha \) is independent of the absolute magnitude of the scattering coefficient, i.e. the concentration of the particles. But it is different for different types of aerosols and mainly influenced by the size distribution of the particles. For a power law number size distribution given by \( dn/dr = n_0 \cdot (r/r_0)^{-\nu} \), the exponent \( \alpha \) can be obtained by \( \alpha = 3 - \nu \) (Junge, 1963). For an aerosol with a size distribution with \( \nu > 3 \), the smaller particles dominate and \( \alpha < 0 \); i.e. the scattering coefficient decreases with increasing wavelength, which is the normal case. If the particles are larger than a few micrometres, then \( \nu < 3 \) and the Ångström exponent is zero or even positive. A strong wavelength dependence of the scattering coefficient (\( \alpha < -1 \)) is an indication of the non-Sahara aerosol, whereas little or almost no dependence on wavelength is a sign of an aerosol containing larger particles, in this study mainly desert particles.

3 Instrument and methods

The volume scattering function of the aerosol has been determined by a custom-made polar nephelometer; its design is similar to the one of Waldram (1945). See Fig. 2: light from a 532 nm solid-state laser of 10 mW power shines into a light trap and illuminates the particles within the beam. A photomultiplier with a collimation optic is mounted on a goniometer and can measure scattered light between 5 and 175°. The scattering volume is approximately 70 mm3 at a 90° scattering angle, increasing to 800 mm3 at 5 and 175° respectively. Measurements are taken at intervals of 5°, in the near-forward direction at 1 to 2° intervals. For a one-volume scattering function a scan from 175 to 5° and back is made and the average is used. One full scan takes 35 min, which can pose a problem with rapidly changing aerosols. Therefore a check for differences between the forward and back scans was made and if the difference was larger than a factor of 1.5, the data were not used. The instrument is enclosed in an airtight housing; by sucking air out of the enclosure, air from outside can be brought into the instrument. Calibration is done by filling the instrument with carbon dioxide, whose volume scattering function is well known. For all data reported below, the scattering of the air molecules has been subtracted; i.e. all scattering functions, phase functions, asymmetry parameters, or backscattered fractions are for aerosol particles only.

Experimentally it is impossible to measure the volume scattering function from 0 to 5°, and 175 to 180°, but the contribution of this range cannot be neglected, especially the forward region. The measured volume scattering function and its shape between 10 and 5° and 170 and 175° can be used to extrapolate the missing regions. This can be done with an accuracy of better than 5% (Horvath, 2015), and thus the complete scattering function is available.

The flow rate through the instrument was \( 3.3 \times 10^{-4} \) m³ s⁻¹. The connection to the sampling inlet of the field laboratory was by a slightly downwards inclined hose of a diameter of 10 mm and a length of 2.5 m. Using data given by Hinds (1999, chap. 10), the loss of particles due to sedimentation in the tube amounted to 10%, 2%, and 0.3% for particles with diameters of 10, 5, and 2 µm. Losses due to dif-

![Figure 2. Principle of the polar nephelometer.](Image 337x549 to 535x727)
to September 2016 there were a total of 15 Sahara dust events, on 96 days out of 153. For details, see https://www.mapama.gob.es/es/calidad-y-evaluacion-ambiental/temas/atmosfera-y-calidad-del-aire/episodiosnaturales2016_tcm30-379284.pdf (last access: 8 December 2018). In June 2016 the events were on the following days: 2–3, 6–11, and 21–30. To distinguish dust aerosols from others, two methods have been applied. (1) The size of the dust particles is larger than 1 µm; thus, the Angström exponent normally is larger than −1, whereas the non-Sahara aerosol has Angström exponents below −1. (2) Using 72 h back trajectories (from the NOAA HYSPLIT website, Draxler and Rolph, 2003), the likely origin of the particles can be estimated. Since the Sierra Nevada is a small mountain range, most of the time the air masses reached the measuring location (at an elevation of 2500 m a.s.l.) without admixing of particles of possible nearby sources. The origin of the air mass was classified in a total of six groups. Figure 3 shows the typical situations: Sahara, Sahara high, Atlantic, Atlantic North, Atlantic/Sahara, North Africa/Mediterranean. From the source region in the Sahara to the receptor region in the Sierra Nevada, the particles travelled around 1500 km. Table 1 characterizes the six types which occurred during this campaign.

A plot of the Angström exponent as a function of time measured with the three-wavelength Integrating Nephelometer is shown in Fig. 4. The classification using the back trajectories is given above. It is evident that for aerosols influenced by the Sahara, the Angström exponent is larger than −1; thus, the two methods of determining Sahara aerosols agree in most of the cases.

### Results

An overview of the scattering coefficient obtained by integration of the measured volume scattering function and the value measured with the Integrating Nephelometer is shown in Fig. 5, and only the light scattering of the particles is plotted; i.e. the Rayleigh scattering of the air was subtracted. The data obtained with the Integrating Nephelometer are values, taken at intervals of minutes, and depict also the variability of the layered aerosol. The scattering coefficient obtained by integration of the polar nephelometer data are averages.
Figure 3. Back trajectories used for classification of aerosol types.

Figure 4. Angström exponent of the spectral scattering coefficient of the aerosol measured and classification by back trajectories. For the desert aerosol the Angström exponent is larger than $-1$. In addition, the asymmetry parameter is plotted too; for desert aerosol it is larger than 0.65.

Figure 5. Measured scattering coefficient of the aerosol during the three periods of observation. The solid red line is the signal of the Integrating Nephelometer; the points are scattering coefficients obtained by integrating the measured volume scattering function.
of about 35 min. Agreement between the two datasets is evident. Three periods of measurements can be seen. During Period 1 mostly a distinct intrusion of Sahara aerosol was observed, except for the beginning; in Period 2 mainly aerosol from the Atlantic reached the site, whereas during Period 3 again Sahara particles dominated the aerosol. The air masses passing over the Atlantic had a much lower scattering coefficient, being about twice the one of pure air.

A plot of all Sahara and non-Sahara phase functions is depicted in Fig. 6. There is some scatter in the data due to the layered aerosol, but it is evident that the phase functions of the Sahara and non-Sahara aerosol are different. The averages of all phase functions of definite Sahara origin and of definite non-Sahara origin and the standard deviations are shown in Fig. 7.

5 Discussion

The comparison between the two types of average phase functions of Fig. 7 shows definite differences. The non-Sahara (mainly Atlantic) phase function has less forward scatter (on average $31\text{ sr}^{-1}$ at $0^\circ$) than the Sahara phase function ($62\text{ sr}^{-1}$). This is readily explained by the larger size of the desert particles. Whereas spherical particles are a good approximation for the near-forward scattering, the smaller backscattering of the Sahara phase function compared to the non-Sahara particles only can be explained by the irregular shape of the particles, for spherical large-particle interferences and resonances are most pronounced, which leads to a considerable increase in backscattering and a low side scattering, both of which have not been observed. For irregularly shaped particles which furthermore are randomly oriented by Brownian rotation, the backscattering is by far less, up to a factor of 10 compared to spherical particles; see e.g. Von Hoyningen-Huene and Posse (1997) or Mishchenko et al. (1997), Mishchenko (2000), and Nousiainen and Kandler (2015).

The complete volume scattering function or phase function can be used to determine derived properties. The scattering coefficient is obtained by integrating the volume scattering function over the full solid angle. A comparison with the calculated and measured scattering coefficients was shown in Fig. 5; the agreement is evident.

For modelling of e.g. radiative transfer, the asymmetry parameter of an aerosol is an important parameter. It is obtained by folding the phase function with the cosine of the scattering angle. For the average phase functions shown in Fig. 7 the asymmetry parameter for the Sahara aerosol particles is 0.71 with a standard deviation of 0.03; for the non-Sahara particles it is $0.56 \pm 0.04$. The difference is significant; there-
Therefore, the asymmetry parameter is a good indicator for the desert aerosol. In Fig. 4 the measured asymmetry parameter is added to the graph of the Angström exponent and the origin of the aerosol. Whenever the origin of the aerosol indicates desert particles, the asymmetry parameter is high. At the same time the Angström exponent also is high. Figure 8 is a plot of all the data points and shows the relationship between the Angström exponent and the asymmetry parameter for all measured phase functions for which Integrating Nephelometer data and polar nephelometer data were available simultaneously. Evidently a larger asymmetry parameter is associated with a larger Angström exponent, but the relationship is not very pronounced.

The scattering function of the desert aerosol has a low back scattering, which is typical for non-spherical particles. Therefore it is to be expected that a characterization of desert aerosol particles could be achieved by considering the fraction of backscattered light. It is defined as the ratio of the integral of the volume scattering function between 90 and 180° divided by the integral over the full angle and is readily available once the volume scattering function is known. A time series of the backscattered fraction obtained in this way is shown in Fig. 9, black squares; the backscattered fraction obtained from polar nephelometer measurements is lower for the aerosol dominated by desert particles, as expected, but less distinct than the asymmetry parameter. Furthermore the backscattered fraction obtained with the Integrating Nephelometer (red line) is systematically larger, which has the following reason: the backscattered fraction for the Integrating Nephelometer is obtained by dividing the signal BbsG (signal of the Nephelometer in the backscattering range for green light) by BsG (signal of the total scattered light). Both signals are proportional to the light flux scattered by the aerosol, but they are truncated, since it is experimentally impossible to integrate the scattered light flux from 90 to 180° and 0 to 180° respectively. But the measured volume scattering functions permits the simulation of the truncation effect; and for a range of the scattering angles between 8 and 170° the simulated BbsG is obtained by integrating the measured volume scattering function from 90 to 170° and BsG is obtained by integrating the measured volume scattering function from 8 to 170°. The simulated truncated backscatter fraction is calculated by dividing the simulated BbsG by the simulated BsG. The expected signal is also shown (blue circles), which is in much better agreement with the Integrating Nephelometer data.
Table 2. Characteristics of the two types of aerosols.

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<th>Sahara aerosol</th>
<th>Non-Sahara aerosol</th>
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<tbody>
<tr>
<td></td>
<td>Value</td>
<td>SD</td>
</tr>
<tr>
<td>Asymmetry parameter (–)</td>
<td>0.71</td>
<td>0.03</td>
</tr>
<tr>
<td>Backscattered fraction (–)</td>
<td>0.094</td>
<td>0.014</td>
</tr>
<tr>
<td>Average scattering coefficient (Mm⁻¹)</td>
<td>71</td>
<td>22</td>
</tr>
<tr>
<td>Phase function at 0° (sr⁻¹)</td>
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<td>27</td>
</tr>
<tr>
<td>Phase function at 90° (sr⁻¹)</td>
<td>0.26</td>
<td>0.03</td>
</tr>
<tr>
<td>Phase function at 180° (sr⁻¹)</td>
<td>0.31</td>
<td>0.11</td>
</tr>
<tr>
<td>Lidar ratio (sr⁻¹)</td>
<td>44</td>
<td>16</td>
</tr>
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</table>

For the Sahara-particle-dominated aerosol the asymmetry parameter is larger and the backscattered fraction is smaller than for the non-Sahara aerosol. So it is obvious to use both parameters to characterize the Sahara aerosol. A plot of all data points in the (b, g) plane is shown in Fig. 10. The points representing the Sahara aerosol and the non-Sahara aerosol are well separated. The curve gives the relationship between the backscattered fraction and the asymmetry parameter calculated for monomodal spherical particles. For other refractive indices or ellipsoidal particles an almost identical curve is obtained; the point of inflection is at a slightly different location. For bimodal size distributions the data points are located to the right of the curve (Horvath et al., 2016), as is the case for the data of this study. A clear distinction between the points representing Sahara and non-Sahara particles is possible. Additionally, points are plotted, which were obtained in Kyoto, Japan, during an event of long-range transport dust intrusion from the Gobi. These data perfectly fit together with the Sahara data.

6 Conclusion

The volume scattering function of the atmospheric aerosol was measured in the Sierra Nevada, where intrusions of Sahara aerosol are frequent. The origin of the aerosol particles was determined by back trajectories and/or by the Angström exponent of the wavelength dependence of the scattering coefficient of the aerosol. Usually it took 48 h or more from the origin to the measuring site. Thus particles larger than 8 µm have been lost by sedimentation and the remaining coarse dust particles were reduced considerably.

The phase function of the Sahara aerosol has more forward scatter, and less back scatter; it is more asymmetric than the non-Sahara aerosol, which in this study was mainly marine, with little continental influence. A few parameters of the aerosol are listed in Table 2.

The best distinction between Sahara and non-Sahara aerosol is possible when using the asymmetry parameter, as is evident also in Figs. 4 and 10. This study suggests that particles causing an asymmetry parameter of the phase function above a value of 0.65 could be considered of Sahara desert origin. The asymmetric scattering of the Gobi desert aerosol is very similar to the one of the Sahara aerosol.

The asymmetry parameter can be used as a distinction between aerosol dominated by desert particles and the other aerosol particles. In this work the desert-dominated particles have an average asymmetry parameter of g = 0.71, whereas the mainly marine aerosol has g = 0.56. The asymmetry parameter of the aerosol measured in Vienna is below 0.65. This aerosol is little influenced by maritime particles and not at all by particles of desert origin; it could be considered as continental background plus traffic-related particles (Horvath et al., 2016). The PM10 is regulated in Europe’s air quality standards to a 24 h average of 50 µg m⁻³ with 35 permitted exceedances per year, and a yearly average of 40 µg m⁻³. Frequent dust intrusions might make it impossible to comply with the regulations; e.g. in 2016 Sahara dust intrusions in the vicinity of the Sierra Nevada were on 96 days. But natural sources are allowed to be excluded from the rules governing the air quality standards if their origin can definitely be proven. In that case the measurement of the asymmetry parameter, possibly in combination with the backscattered fraction, definitely will create clarity.

Data availability. All measured phase functions are available in the Supplement.

Supplement. The supplement related to this article is available online at: https://doi.org/10.5194/acp-18-17735-2018-supplement.

Author contributions. LAA and FJOR planned the SCOPE project and its financing, supplied the infrastructure in the Sierra Nevada site, took care of the operation and data analysis of the Integrating Nephelometer as well as many other instruments, and supplied data on dust events. HH designed and operated the polar nephelometer, reduced the data and performed the data analysis, wrote the draft and after through discussion with all authors finalized the paper.

Competing interests. The authors declare that they have no conflict of interest.

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References


