Tropospheric ozone variability in the tropics from ENSO to MJO and shorter timescales

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Abstract. Aura OMI and MLS measurements are combined to produce daily maps of tropospheric ozone beginning October 2004. We show that El Niño-Southern Oscillation (ENSO) related inter-annual change in tropospheric ozone in the tropics is small in relation to combined intra-seasonal/Madden–Julian Oscillation (MJO) and shorter timescale variability by a factor of ∼3–10 (largest in the Atlantic). Outgoing longwave radiation (OLR), taken as a proxy for convection, suggests that convection is a dominant driver of large-scale variability of tropospheric ozone in the Pacific from inter-annual (e.g., ENSO) to weekly periods. We compare tropospheric ozone and OLR satellite observations with two simulations: (1) the Goddard Earth Observing System (GEOS) chemistry-climate model (CCM) that uses observed sea surface temperatures and is otherwise free-running, and (2) the NASA Global Modeling Initiative (GMI) chemical transport model (CTM) that is driven by Modern Era Retrospective-Analysis for Research and Applications (MERRA) analyses. It is shown that the CTM-simulated ozone accurately matches measurements for timescales from ENSO to intra-seasonal/MJO and even 1–2-week periods. The CCM simulation reproduces ENSO variability but not shorter timescales. These analyses suggest that a model used to delineate temporal and/or spatial properties of tropospheric ozone and convection in the tropics must reproduce both ENSO and non-ENSO variability.

1 Introduction

The El Niño-Southern Oscillation (ENSO) and its effects on the atmosphere and ocean have been extensively studied and documented. Trenberth (1997) provides several key references with an overview description and historical account of ENSO. The terminology, ENSO, is understood to consist of El Niño (warmer than average ocean temperatures in the tropical eastern Pacific – i.e., warm phase) typically followed by La Niña (cooler than average ocean temperatures in the tropical eastern Pacific – i.e., cool phase). ENSO events have ∼2–7-year timescales and produce planetary-scale changes in tropical sea surface temperature (SST), convection, and winds. Peak activity for ENSO occurs generally centered about Northern Hemisphere autumn to mid-winter months (e.g., largely October–January).

The effects of El Niño on atmospheric composition, including ozone, have been studied from both satellite and ground-based measurements (e.g., Chandra et al., 1998; Fujiwara et al., 1999; Thompson et al., 2001; Nasser et al., 2009; Lee et al., 2010; Ziemke et al., 2010; Randel and Thompson, 2011; Neu et al., 2014), global chemical transport models driven by specified meteorology (e.g., Valks et al., 2003; Duncan et al., 2003; Chandra et al., 2009; Murray et al., 2013) and general circulation models (GCMs) (e.g., Sudo and Takahashi, 2001; Zeng and Pyle, 2005; Doherty et al., 2006; Randel et al., 2009; Oman et al., 2011; Sekiya and Sudo, 2014). Tropospheric ozone is important as both a greenhouse gas and precursor of the hydroxyl radical (OH), the primary atmospheric oxidant. Tropospheric ozone in the tropics is especially sensitive to changes in deep convection associated with ENSO. An increase (decrease) in dynami-
cal convection from ENSO events in the tropical Pacific induces a decrease (increase) in tropospheric column ozone. Although changes in convection are fundamentally dynamical there are also ENSO-related changes in composition that affect ozone precursors, such as increases in emissions from biomass burning over Indonesia due to suppressed rainfall during El Niño. There can also be long-range transport effects on tropospheric ozone at northern mid-latitudes related to ENSO including induced trends over long records. Lin et al. (2014) studied the effects of ENSO/Pacific Decadal Oscillation (PDO) on tropospheric ozone at Mauna Loa Observatory (19.5°N, 156.5°W, altitude 3.4 km). By combining 40 years of surface ozone measurements with a set of chemistry-climate model simulations they found that the flow of ozone-rich air from Eurasia towards Hawaii during spring weakened in the 2000s as a result of La Niña-like decadal cooling in the tropical Pacific. This circulation-driven ozone decrease offsets the ozone increase that otherwise would have occurred at Mauna Loa in spring due to rising Asian anthropogenic emissions.

Ziemke et al. (2010) produced a monthly tropospheric ozone ENSO index (OEI) over a multi-decadal time record (beginning 1979) by differencing satellite-measured column ozone in the tropics between the eastern and western Pacific. They noted that the OEI could be used as a diagnostic test for modeled ozone including tropospheric ozone sensitivity relating to changes in SSTs. Oman et al. (2011) found excellent agreement between the measured OEI with the OEI produced by the Goddard Earth Observing System (GEOS) free-running chemistry-climate model (CCM) with observed SSTs over a 25-year period. This demonstrated an appropriate response of the CCM meteorology to the ENSO signature of the imposed SSTs; the fidelity of the ozone response to the induced circulation and photochemical changes included realistic horizontal and vertical gradients in tropospheric ozone.

The tropical atmosphere and ocean exhibits intra-seasonal and shorter timescale variability with periods much shorter than ENSO from days or weeks to several months. The leading source of intra-seasonal variability is related to the Madden–Julian Oscillation (MJO) with characteristic timescales of about 1–2 months (Madden and Julian, 1971, 1994). The MJO is identified as large-scale circulation cells in equatorial latitudes that propagate eastward from the Indian Ocean to at least the central Pacific. In the original discovery of the MJO, Madden and Julian (1971) described the oscillation as a 40–50-day variation in surface pressure, zonal winds and temperature at different levels of the troposphere. Madden and Julian (1994) note that zonal wind anomalies in the upper troposphere associated with the MJO sometimes traverse the entire circumference of the Earth. The strongest variability for the MJO occurs around northern wintertime months when the intensity of ENSO events is largest. The MJO modulates regional monsoon (in particular the Australian and Indian monsoon) which impacts particulate matter (e.g., Ragsdale et al., 2013) and surface ozone (e.g., Barrett, et al., 2012), both of which contribute to poor air quality in the tropics and/or subtropics. The ocean-atmosphere coupling associated with the MJO may also affect the duration and onset of ENSO (e.g., Hoell et al., 2014, and references therein). The MJO alters stratospheric circulation including the strength of the Northern Hemisphere polar vortex and timing of stratospheric sudden warming events (e.g., Garfinkel et al., 2012, 2014, and references therein) and also modulates tropical Kelvin waves (Guo et al., 2014). Wheeler and Hendon (2004) quantify the MJO using two leading derived Empirical Orthogonal Functions (EOFs) of combined tropospheric zonal winds and Outgoing Longwave Radiation (OLR) in the tropics. Their method quantifies the MJO into eight separate identifiable temporal phases beginning from onset extending through the full 1–2 month cycle.

Using a chemical transport model (CTM) and measurements from the Aura Tropospheric Emission Spectrometer (TES), Sun et al. (2014) indicated that the MJO in tropospheric ozone in tropical latitudes may be locally up to 47 % of total variability. Their estimate is comparable to the ~5–10 Dobson units (DU; 1 DU = 2.69 × 10^20 molecules m^-2) MJO variability (out of ~15–20 DU background ozone) in troposphere ozone in the tropical Pacific by Ziemke and Chandra (2003). One of the results of Sun et al. (2014) was that large-scale advection within the CTM explains most of the simulated changes in ozone relating to the MJO.

In addition to ENSO and intra-seasonal/MJO changes, Dunkerton and Crum (1995) showed that there is considerable convective variability in the tropics with shorter timescales of 2–15 days. Dunkerton and Crum (1995) used daily outgoing longwave radiation (OLR) in the tropics to relate 2–15-day disturbances with intra-seasonal oscillations/MJO signals and found distinction between them as well as moderate interaction between them during convectively active phases of the intra-seasonal oscillations. A long existing problem with GCM/CCM simulations is difficult in producing a realistic MJO in the atmosphere (e.g., Lin et al., 2006; Hung et al., 2013; Del Genio et al., 2015; and references therein). Efforts have demonstrated that there is a causal link between how well gross moist stability and vertical advection is treated in models with how well those models reproduce a variation similar to the MJO (e.g., Benedict, et al., 2014, and references therein).

The purpose of our study is to characterize the variability of tropical tropospheric ozone for timescales ranging from ENSO to MJO and shorter time periods in relation to tropical convection and atmospheric model simulations of ozone. We compare observed tropospheric ozone with ozone simulated from two NASA Goddard models of atmospheric composition, one being a CCM forced by observed monthly SSTs and the other a CTM driven by meteorological reanalyses. Section 2 discusses data and models for our analysis while Sect. 3 describes the impact of ENSO- vs. non-ENSO-related changes in tropospheric ozone in relation with con-
vective forcing. Section 4 describes derivation of a useful tropospheric ozone diagnostic from OMI/MLS while Sect. 5 shows some of its applications as applied to model ozone and OLR measurements. Finally, Sect 6 provides a summary.

2 Data and models

Daily measurements of tropospheric column ozone (TCO) in tropical latitudes are calculated using the OMI/MLS residual method of Ziemke et al. (2006). This method subtracts MLS stratospheric column ozone from OMI total column ozone for near clear-sky scenes (i.e., radiative cloud fractions < 30%). Ziemke et al. (2014) evaluated three other OMI/MLS TCO products and concluded that the Global Modeling and Assimilation Office (GMAO) data assimilation product was best to use overall when considering all factors including global coverage and ozone profile information. However, Fig. 12 of Ziemke et al. (2014) showed that the assimilation product when limited to tropical latitudes had zonal variability ~ 10–15 DU in stratospheric column ozone which was considerably larger than direct satellite measurements that typically have zonal variability of only a few DU. In addition, this larger zonal variability in stratospheric column ozone coincided with a reduced zonal wave-one pattern of TCO with assimilation, also considered inconsistent with previous TCO measurements. Our main reason for using the product of Ziemke et al. (2006) for the tropics stems from being independent of MERRA/GEOS-5 analyses including winds used by both the assimilation and trajectory ozone products. There are known errors with tropical stratospheric winds in the analyses caused by spurious transport (Tan et al., 2004). Although Tan et al. (2004) diagnosed an older assimilation system, comparisons of MLS ozone and N_2O with our CTM simulations using MERRA meteorology show that spurious transport in the tropical and subtropical lower stratosphere is still a problem. These errors in stratospheric winds with assimilation produce errors in the derived ozone profiles including TCO.

The OMI/MLS residual product combines MLS v3.3 ozone profiles with OMI version 8.5 total ozone measurements. Data quality and description of the MLS v3.3 ozone profiles are discussed by Livesey et al. (2011). Description and access to the OMI data may be obtained from the NASA webpage http://disc.sci.gsfc.nasa.gov/Aura/data-holdings/OMI. Horizontal gridding for TCO is 1° latitude × 1.25° longitude. The OMI/MLS residual ozone product uses WMO NCEP 2 K km⁻¹ lapse-rate tropopause pressure to separate tropospheric from stratospheric ozone. Our study also uses OLR daily measurements for 2004–2012 at 2.5° × 2.5° horizontal gridding obtained from the National Oceanic and Atmospheric Administration (NOAA) webpage http://www.esrl.noaa.gov/psd/data/gridded/. OLR is the amount of radiative flux (units W m⁻²) re-emitted back to space in the 3.55–3.91 µm wavelength band.

The Global Modeling Initiative (GMI) CTM hindcast simulation includes a chemical mechanism suitable for the troposphere and stratosphere (Duncan et al., 2008; Strahan et al., 2007). Although the CTM simulation extends from 1979 through 2012 we include a limited period of October 2004 through December 2012 to coincide with the OMI/MLS measurements. The emissions of trace gases and aerosol fields used in the CTM simulations are described by Duncan et al. (2008), however, anthropogenic emissions have been updated and include year-specific scaling factors (van Donkelaar et al., 2008). Anthropogenic and biomass global emissions include surface emissions from industry/fossil fuel, biomass burning, biofuel combustions, and contributions from aircraft. Biomass burning emissions in the CTM are from van der Werf et al. (2010) and are extended through year 2010. Observationally-based biomass burning emissions are used in the CTM through year 2010 with the 2010 emissions repeated for 2011–2012. Most of the global emissions such as fossil fuel, biofuel, and biomass burning emissions for the CTM represent monthly means; however, lightning NO_x and biogenic emissions (such as isoprene) are calculated online within the model and can vary daily. More detailed description of emissions for this simulation is given by Strode et al. (2015). The CTM meteorological fields are taken from Modern-Era Retrospective Analysis for Research and Applications (MERRA) (Rienecker et al., 2011).

The CCM is described by Oman et al. (2013). This CCM is forced by monthly SSTs and specified boundary conditions and fluxes of important greenhouse gases including carbon dioxide, methane, and nitrous oxide. The CCM uses observed monthly mean SSTs over the 1960–2012 time record (Rayner et al., 2003). All global emissions for the CCM including biomass burning and non-biomass burning/anthropogenic were chosen to closely match emissions for the CTM. Lightning NO_x for the CCM varies daily as with the CTM. We again refer to Strode et al. (2015, and references therein) for quantitative details regarding emissions. Both the CCM and CTM tropopause pressure use the WMO lapse-rate definition.

3 ENSO vs. non-ENSO timescale changes in tropical tropospheric ozone

Variability of tropospheric ozone from OMI/MLS and the CTM, shown in Fig. 1 for ENSO, non-ENSO, and intra-seasonal oscillation (ISO) timescale changes, is derived from calculated standard deviation (see figure caption). Also plotted in Fig. 1 are corresponding calculations for OLR, scaled by a factor of 0.18 for plotting with ozone. Figure 1 is comprised of two sets of calculations. Figure 1a corresponds to variability calculated using original daily time series whereas Fig. 1b is the same but with all daily time series filtered to include only extreme ENSO events. For the Niño 3.4 index, ENSO events as defined by NOAA occur when Niño 3.4 is
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Figure 1. (a) Variability in deseasonalized OMI/MLS daily tropospheric column ozone (solid curves), GMI CTM (dotted curves), and OLR (long dashed curves) for ENSO signal, intra-seasonal oscillation (ISO) signals, and with ENSO signals removed (non-ENSO). ISO curves involved band-pass filtering the time series for 25–65-day periods. OLR (units W m\(^{-2}\)) was multiplied by a factor of 0.18 for plotting with ozone. The plotted variability was calculated using amplitude of 2\(\sigma\) to estimate peak-to-peak change. The time record is 1 October 2004–31 December 2012 and all original time series were averaged over 20\(^\circ\)S–20\(^\circ\)N. The ENSO signals were extracted using the linear regression \(T(t) = \beta \times \text{Nino}34(t) + \epsilon(t)\), where \(T\) is original gridded time series, \(t\) is day index, \(\beta\) is a derived constant, \(\text{Nino}34(t)\) is the Niño 3.4 ENSO index, and \(\epsilon(t)\) is the residual (i.e., \(\epsilon(t)\) is identically the non-ENSO component of the time series). A daily Nino34(t) time series was generated from the NOAA monthly record using linear extrapolation. All line curves in Fig. 1 represent 20\(^\circ\)S–20\(^\circ\)N averages as function of longitude. The ISO variability (middle curves) involved band-pass filtering of the original time series for 25–65-day periods (see Appendix A and Fig. 1 caption).

In Fig. 1a ENSO contributes a relatively small amount to the total daily variability of tropospheric ozone in the tropics at all longitudes. Figure 1b shows that this is the case even when only extreme ENSO events are considered, although for extreme events the ENSO variability increases in the Pacific relative to shorter timescales. In Fig. 1a and b ENSO-related change in tropospheric ozone and OLR (bottom curves) is generally smaller than either non-ENSO change (top curves) or ISO timescale changes (middle curves). The CTM reproduces all of the OMI/MLS tropospheric ozone zonal patterns for all three timescale scenarios. Most of the non-ENSO related changes involve the intra-seasonal/MJO and shorter time period changes. These changes are larger than ENSO by a factor of \(\sim 3–4\) in the Pacific and a factor of 10 or more in the Atlantic in both Fig. 1a and b.

4 The daily ozone dipole index (ODI)

We calculate a quantity that we refer to as the ozone dipole index (ODI). This differs from the monthly OEI used by Ziemke et al. (2010) in that it is calculated using daily measurements rather than monthly means and does not include the final 3-month running average that is applied to the OEI. We use this ODI as a diagnostic test for evaluating OMI/MLS tropospheric ozone with other atmospheric parameters, including satellite-measured OLR and similar tropospheric ozone derived from models. The ODI is the deseasonalized difference of western minus eastern Pacific TCO time series each day over the Aura record. Deseasonalization of time series is explained in Appendix A. The ODI calculation involves first averaging TCO from OMI/MLS each day in the tropics over the broad eastern and western Pacific regions (i.e., 15\(^\circ\)S–15\(^\circ\)N, 110–180\(^\circ\)W and 15\(^\circ\)S–15\(^\circ\)N, 70–140\(^\circ\)E, respectively) followed by computing the difference of western minus eastern Pacific. As with the monthly OEI, this differing removes measurement offsets or drifts with time that would be common to both Pacific time series. We also calculate a daily dipole index time series for National Oceanic and Atmospheric Administration (NOAA) OLR measurements in the exact same manner as calculation of the ODI for investigating connections between tropospheric ozone and convection in the Pacific.

either greater than 0.5 (El Niño) or less than –0.5 (La Niña) for five consecutive months. Extreme ENSO events in Fig. 1b were subjectively chosen here to correspond with Niño 3.4 index being either greater than 1.0 or less than –1.0.

ENSO signals (bottom curves) in Fig. 1 were extracted at each grid point from original deseasonalized ozone and OLR time series using the linear regression \(T(t) = \beta \times \text{Nino}34(t) + \epsilon(t)\), where \(T\) is original gridded time series, \(t\) is day index, \(\beta\) is a derived constant, \(\text{Nino}34(t)\) is the Niño 3.4 ENSO index, and \(\epsilon(t)\) is the residual (i.e., \(\epsilon(t)\) is identically the non-ENSO component of the time series).
Statistical coherence and phase of coherence are calculated between the measured ODI and the ODIs derived from the CTM and CCM. These statistics are also calculated between the measured ODI and the OLR daily dipole series. Coherence, a normalized statistic with values lying between 0.0 and 1.0, provides evaluation of statistical connection between two time series as an explicit function of frequency. We refer the reader to Appendix A for details regarding these calculations.

5 Comparisons between measured and modeled ODI

In Fig. 2a we compare time series of measured ODI (red curve) and CTM ODI (dotted blue curve). The two time series appear remarkably similar for timescales varying from low-frequency ENSO to 1–2-month periods (e.g., MJO) and even shorter. Figure 2b is the same as Fig. 2a but for the CCM instead of CTM. The CCM in Fig. 2b reproduces ENSO variability and appears to produce variability at shorter timescales similar to the CTM; however, the evaluation of the models requires more than just visual inspection of time series.

We calculate coherence and coherence phase as functions of frequency to establish a statistical connection between measured and simulated ODIs on varying timescales. The coherence and coherence phase calculated between the OMI/MLS and CTM ODIs are shown in Fig. 3a where square of coherence is shown in the top panel with coherence phase on the bottom. Time periods in days are printed along the horizontal frequency axes for all panels in Fig. 3.

If a simulated ODI exactly matched that obtained from OMI/MLS then the squared coherence would be 1.0 and the phase shift would be 0.0 over the entire frequency spectrum. For the CTM in Fig. 3a, statistical significance of squared coherence exceeds the 99% level for values greater than 0.684. The CTM squared coherence exceeds this value for a broad range of timescales from ENSO (at far left in panel) to the MJO (30–60 days), down to timescales as short as 7–14 days. The excellent agreement in Fig. 3a over broad timescales attests to the realism of the input meteorology and computed photochemistry within the CTM. Figure 3b shows similar calculations for the CCM. The squared coherence in Fig. 3b (top) is statistically significant for ENSO but not shorter timescales. In addition the phase between OMI/MLS and the CCM in Fig. 3b (bottom) is near zero only for very low-frequency ENSO variability.

Convection activity is inferred using OLR flux measured from NOAA polar orbiting satellites (e.g., Chelliah and Arkin, 1992; Liebmann and Smith, 1996). Clouds that are high in the troposphere have cloud-top temperatures colder than cloud tops lying below. The colder cloud tops coincide with reduced OLR and therefore low OLR corresponds to deep convection.

Comparison of the OMI/MLS ODI with the OLR dipole series in Fig. 4 indicates that convection is the main driver of the ODI from ENSO to MJO and shorter periods. Aside from convection/advection forcing, the variability of precursors may also affect the variability of tropospheric ozone on different timescales; however, chemical timescale vs. dynamical timescale must be considered. As an example, CO is a precursor of tropospheric ozone with an average lifetime of

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Figure 3. This figure plots calculated coherence and phase of coherence between OMI/MLS ODI and model (i.e., CTM and CCM) ODI as functions of frequency (periods in days shown). (a) Top panel: coherence-squared between OMI/MLS ODI and CTM ODI. Included are confidence levels for coherence-squared of 95% (i.e., value of 0.393), 99% (value of 0.536), and 99.9% (value of 0.684). Bottom panel: phase of coherence in degrees. Panel (b) is the same as (a) but for the CCM instead of CTM (see Appendix A for details of these calculations).

∼ 2 months (e.g., Petrenko et al., 2013). Conversion of CO to ozone will have a relatively long timescale compared to daily or weekly variability, but not when compared to intra-seasonal to inter-annual variability. As a test, we repeated our analyses where all emissions for the CTM were held constant in time (figures not shown). We found that the variability of CTM ozone such as that shown in Fig. 1a and the coherence/phase in Fig. 3a were nearly identical for the constant-emissions simulation. This suggests that the variability of precursors is not important overall in affecting tropospheric ozone variability on these timescales and on planetary scales. However, the variability of precursors on regional scales can be significant. It was shown by Ziemke et al. (2009) using the CTM and OMI/MLS ozone that biomass burning over Africa, South America, and Indonesia can generate 10–25% and even greater increases of tropospheric ozone in localized regions within or near the burning. The high coherence calculated between measured ODI and the OLR dipole series from inter-annual (i.e., ENSO) to shorter timescales suggests that convection has a dominant influence in forcing large-scale changes in tropospheric ozone in the tropical Pacific. The behavior of OLR with ozone in Figs. 1 and 4 indicates further that convection in the MERRA analyses is being well simulated from ENSO down to weekly timescales.

Figure 5 compares calculated spectral amplitudes of the ODI obtained from OMI/MLS data (red curve), CTM output
Figure 5. This figure plots calculated spectral amplitudes (in DU) of ODI derived from OMI (solid red curve), CTM (dotted blue curve), and CCM (dashed green curve) as functions of frequency (periods in days shown). Spectral amplitude is defined as the square root of \(c(\omega)^2 + s(\omega)^2\), where \(c\) and \(s\) denote Fourier cosine and sine coefficients, \(\omega\) is frequency and the over-bar denotes application of a smoothed spectral estimator (see Appendix A for details of these calculations).

The spectral amplitudes for OMI/MLS and the CTM in Fig. 5 are everywhere comparable and the variability shown by peaks and valleys as functions of frequency are closely identical for periods even shorter than \(\sim 30\) days. In comparison, the CCM has considerably smaller amplitudes at all frequencies and the frequency variability of spectral amplitudes is not consistent with the observations. The spectral analysis including the coherence/coherence-phase statistics moves beyond visual inspection of time series to give a quantitative measure of model performance.

Power spectra for TCO time series averaged over the Indian Ocean just north of the Equator are shown in Fig. 6 for OMI/MLS ozone and the CTM and CCM simulated ozone. This tropical region is where the 1–2-month MJO signal-to-noise in tropical TCO maximizes for both data and the CTM. MJO variability in Fig. 5 has well-defined peak amplitudes around 45–50-day period for both the data and the CTM. However the CCM power spectrum does not show any consistent MJO or shorter timescale variability and essentially only generates an ENSO variation at very low frequency.

Power spectra for TCO time series averaged over the Indian Ocean just north of the Equator are shown in Fig. 6. All three panels show calculated power spectra (in units of DU^2) of daily tropospheric column ozone time series averaged over a broad region of the tropical Indian Ocean (0–10° N, 70–80° E) where the MJO signal is statistically significant well above 95 % for OMI/MLS and the CTM. The top, middle, and bottom panels are for OMI/MLS data, CTM output, and CCM output, respectively. The power spectra are plotted vs. frequency with periods in days shown. A power spectrum is defined by \(|c(\omega)|^2 + |s(\omega)|^2\)/2 where \(c\) and \(s\) denote derived Fourier cosine and sine coefficients, \(\omega\) is circular frequency and the over-bar denotes application of a smoothed spectral estimator. Estimated background noise is denoted “BG” with 95 % confidence level shown in each panel (see Appendix A for details of these calculations).

6 Summary

We have studied the variability of tropospheric ozone in the tropics from ENSO to intra-seasonal/MJO and weekly timescales using satellite measurements and two simulation models. Aura OMI and MLS satellite measurements are combined to derive daily maps of tropospheric ozone for October 2004 through 2012. Daily OLR from NOAA for the same time record are included to relate tropospheric ozone variability to changes in convection. The two models that we use are (1) the free-running GEOS Chemistry-Climate Model (CCM) and (2) the Global Modeling Initiative (GMI) chemistry-transport model (CTM) driven by Modern-Era Retrospective Analysis for Research and Applications (MERRA) meteorological analyses.

Non-ENSO timescale changes in measured tropospheric ozone and convection in the tropics are found to be larger than ENSO-related changes by a factor of about 3–4 in the Pacific and up to a factor of \(\sim 10\) in the Atlantic. The non-ENSO variability in tropospheric ozone and convection is comprised mostly of intra-seasonal/MJO to 1–2-week timescale changes. Time series analysis including coherence calculations with OLR satellite data suggests that large-scale variability of tropospheric ozone in the Pacific from ENSO to weekly timescales is driven largely by convection.

We developed a tropospheric ozone dipole index (ODI) from OMI/MLS measurements by differencing western mi-

Our analyses show that the Goddard CTM reproduces ozone observations exceptionally well over timescales from ENSO down to weekly periods whereas the Goddard CCM reproduces only ENSO variability. The inability of the CCM to generate shorter timescales such as an MJO is a known problem with GCMs/CCMs. Using daily instead of monthly SSTs would likely not improve performance of the CCM in light of previous studies. Del Genio et al. (2015) suggest that for these models to generate an MJO they need to have cloud/moisture-radiative interactions and convection-moisture sensitivity. Understanding the differences in ozone variability between the CCM and CTM can help quantify possible missing or inaccurate feedback processes as future work. An important result we find is that using a model to quantify temporal and spatial properties of tropospheric ozone in the tropics requires that the model properly simulate the non-ENSO variability which includes the MJO and shorter periods.
Appendix A

A1 Estimated precision errors for OMI/MLS TCO including calculated ODI

Estimated root-mean-square (RMS) precision errors for OMI/MLS 1° × 1.25° daily gridded TCO are given by Kar et al. (2010). Precision values in the extra tropics were shown to be up to ∼10 DU or greater while in tropical latitudes values were smaller at ∼5 DU. Figure A1 shows daily time series of eastern and western Pacific OMI/MLS TCO used to calculate the ODI (the two Pacific regions are defined in the figure caption). The ODI follows by taking the western minus eastern Pacific TCO each day followed by deseasonalizing this difference time series (deseasonalization is discussed in Sect. A2 below). The time series in Fig. A1 appear generally of opposite signature with evidence of some temporal phase shifts for intra-seasonal and shorter timescales. An El Niño (La Niña) condition coincides when these two time series have largest (smallest) separation on inter-annual timescale.

RMS precision errors for the time series in Fig. A1 were obtained by taking local daily RMS uncertainties at 1° × 1.25° resolution and adjusting these numbers by the spatial averaging invoked. By taking an upper bound of 10 DU for this number and dividing it by $\sqrt{N}$ (N is the total number of the grid points included in the spatial averaging) we get an estimate of time series precision. (This precision estimate represents standard error of the mean.) Dividing by $\sqrt{N}$ assumes that tropospheric ozone measurements detected by OMI are stochastically independent. For either the western or eastern Pacific region encompassing a domain of 30° latitude × 70° longitude there are a total of 1680 grid points. Largely because of applied cloud filtering (i.e., cloud fractions < 30 %) the actual average number of grid points is about 680 (i.e., $N = 680$). This yields $10/\sqrt{680} = 0.38$ DU as an estimated precision for either eastern or western Pacific time series in Fig. A1. An estimate of precision for the daily ODI is then $\sqrt{0.38^2} = 0.54$ DU assuming stochastic independence between the two regions.

A2 Spectral analysis

Koopmans (1974) details calculation of coherence and its phase using Fourier spectral analysis with smoothed spectral estimators. All daily ozone time series in our study were deseasonalized prior to any Fourier analysis. Deseasonalization was accomplished by first applying a low-pass filter (with half-amplitude filter response at 60-day period and zero phase shift at all periods) to original daily time series; this was followed by averaging similar days over consecutive years to obtain a 365-day pseudo-climatology for the annual cycle. This estimated annual climatology was then subtracted from original daily time series for each consecutive year. Potential leakage of nearby Fourier cosine and sine coefficients was reduced by applying a tapered cosine window to deseasonalized time series with 25 % cosine tapering at each end (e.g., Harris, 1978). For all derived spectra including cross-spectra for coherence we applied a Daniell seven-point smoothed spectral estimator. Resulting critical coherence at 95, 99, and 99.9 % confidence levels is 0.627, 0.732, and 0.827, respectively.

Power spectra with estimated 1–2-month signal-to-noise were calculated in the tropics for OMI/MLS and CTM TCO similar to Ziemke et al. (2007). Figure 6 in the main text shows power spectra with estimated signal-to-noise for both background and 1–2 month signal for the Indian Ocean region where the MJO signal for both OMI/MLS and CTM TCO is largest. In Fig. 6, an estimated background noise power spectrum (i.e., denoted BG) for each time series was estimated using a first-order autoregressive model \( T(t) = \alpha \times T(t - 1) + N(t) \), where \( \alpha \) is a derived constant, \( t \) is the day index, and \( N(t) \) is normally distributed random noise with mean of zero. For power spectra using the seven-point estimator the 95% critical signal-to-noise ratio level is 1.69.

A3 ENSO vs. non-ENSO variability

The top panel in Fig. A2 shows OMI/MLS time series for the ENSO component (thick curve), non-ENSO component (thin curve), and annual cycle (dotted curve) in the tropical western Pacific. The bottom panel in Fig. A2 is the same as the top panel but instead for the CTM. The selected region for these time series is 10–20° S, 115–120° E which coincides with largest ENSO variability for both OMI/MLS and CTM TCO. The ENSO variability was extracted using linear regression (see figure caption). Figure A2 shows that the CTM closely tracks OMI/MLS measurements for the non-ENSO components. ENSO variability for both the CTM and OMI/MLS is smaller than non-ENSO (comprised mostly of MJO and shorter timescales).
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