First detection of tidal behaviour in polar mesospheric water vapour by ground based microwave spectroscopy

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Abstract. Mesospheric water vapour has been observed above ALOMAR in northern Norway (69° N 16° E) by our group since 1995 using a 22 GHz ground based microwave spectrometer. A new instrument with higher sensitivity, providing a much better time resolution especially in the upper mesosphere, was installed in May 2008. The time resolution is high enough to provide observations of daily variations in the water vapour mixing ratio. We present the first ground based detections of tidal behaviour in the polar middle atmospheric water vapour distribution.

Diurnal and semidiurnal variations of water vapour have been observed and due to the long chemical lifetime of water they are assumed to be caused by changing wind patterns which transport water-rich or poor air into the observed region. The detected tidal behaviour does not follow any single other dynamical field but is instead assumed to be a result of the different wind components.

Both the diurnal and semidiurnal amplitude and phase components are resolved. The former shows a stable seasonal behaviour consistent with earlier observations of wind fields and model calculations, whereas the latter appears more complex and no regular behaviour has so far been observed.

1 Introduction

The dynamical behaviour of the polar middle atmosphere is important for the global circulation and in order to correctly model it a comprehensive knowledge of all the involved processes is needed.
for observations of the water vapour mixing ratio (WVMR) on different time scales, from long term trends and seasonal behaviour to isolated events such as sudden stratospheric warmings (Hartogh and Jarchow, 1995; Nedoluha et al., 1996; Seele and Hartogh, 2000; Hartogh et al., 2010). The chemical lifetime of water vapour in the middle atmosphere is on the order of days to weeks (depending on altitude) which is similar to the characteristic dynamical time scale of interest and can therefore be used as a tracer molecule (Brasseur and Solomon, 1998). In contrast to tidal behaviour in the dynamical fields, almost no information of the behaviour of atmospheric tracers such as water vapour exist to a large extent due to the observational difficulties. Haefele et al. (2008) report of daily variations in the WVMR measured with a ground based microwave spectrometer in the Alps. They compare these variations to chemistry-climate models (SOCOL, MSDOL and LMDz) and conclude that the variations in the stratosphere are mainly induced by meridional advection and in the mesosphere by vertical advection. At high latitudes such as the polar regions the tidal amplitudes are generally understood to be smaller than at the equator due to the much weaker solar excitation. To the best of the authors knowledge this is the first time diurnal (and semidiurnal) variations of water vapour at a polar location have been observed and described.

2 Instrument

To retrieve the tidal variations in the WVMR in the middle atmosphere we use a ground based microwave spectrometer. Our observations of the polar, middle atmospheric WVMR by ground based spectroscopy at the ALOMAR observatory in northern Norway (69.16° N, 16.00° E) started with the successful WASPAM instrument (Wasserduckmpf- und Spurengasmessung in der Atmosphäre mit Mikrowellen), described in Hartogh and Jarchow (1995). Observations from WASPAM have been used to gain insight in the annual variability of water vapour as well as its behavior during sudden stratospheric warmings and the quasi 5-day wave (Seele and Hartogh, 1999, 2000; Sonnemann et al., 2008; Hartogh et al., 2010). A new instrument, cWASPAM (cooled-WASPAM), was installed at ALOMAR in May 2008 and has been running continuously ever since, with the exception of a few maintenance stops. cWASPAM has a higher sensitivity compared to WASPAM which was achieved by the following improvements: cooling of the horn antenna (which was not done for WASPAM), use of an ortho mode transducer in order to use both vertical and horizontal polarisation in the analysis and InP amplifiers cooled to 15 K instead of GaAs cooled to 20 K. A more extensive description of the instrument can be found in Hallgren et al. (2010) and Hallgren (2010). A sister instrument, cWASPAM3, recently participated in an intercomparison campaign at Zugspitze Observatory where all participating instruments agreed well on their retrieved water vapour profiles (Straub et al., 2011).

The two main objectives of the new instrument were to replace the WASPAM instrument (Hartogh and Jarchow, 1995) installed 1995 and to detect and characterise the short term variations (such as tidal behaviour) of the WVMR in the polar mesosphere. Similar to its predecessor, cWASPAM observes the rotational transition of the water molecule at 22.235 GHz. The vertical and horizontal polarisation components of the signal are simultaneously analysed with two Chip Transform Spectrometer (CTS) backends (Hartogh and Hartmann, 1990; Villanueva and Hartogh, 2004; Villanueva et al., 2006; Paganini and Hartogh, 2009). The CTS are identical and have a bandwidth of 40 MHz which is divided into 4096 channels with a resolution of 14 kHz each. The channels are slightly overlapping so that the effective spectral resolution becomes 10 kHz. Due to pressure broadening of the emission the narrow bandwidth acts as the practical lower boundary of the measurements (approximately 40–45 km). The upper limit is set by the transition altitude where the Doppler broadening becomes more important than the pressure broadening. Typical Doppler broadening of this line is \(\sim 30\) kHz, well above our channel resolution. The transition for this particular line occurs at 80–85 km. Above this altitude we can no longer resolve the vertical distribution of water vapour.

The noise temperature of each polarisation backend is on the order of 30 K. The spectra derived from the vertical and horizontal polarisation are statistically independent so that averaging the two signals essentially doubles the integration time, or reduces the noise by \(1/\sqrt{2}\). Although the backends are identical the output differs slightly due to imperfections in the polarisation splitter (an ortho-mode transducer). Hence they are averaged by the standard way of combining two independent measurements \((x_1\) and \(x_2\)) weighted by their respective spectral noise. This results in the equivalent single polarisation backend with a noise temperature of approximately \(20\) K. Even with a moderate tropospheric transmission of 50\% the high sensitivity of the instrument allows retrieval of reliable atmospheric profiles throughout most of the observable range every six hours. To study the dynamics in the lower range of the sensitive region (e.g. sudden stratospheric warmings) the integration time can be shortened to four hours.

3 Data retrieval

The retrieval follows the optimal estimation method (OEM) described in Rodgers (1976). It requires temperatures of the background atmosphere which are compiled by a combination of actual weather data from NCEP (McPherson et al., 1979) and the CIRA86 model (Fleming et al., 1990). During summer season the CIRA86 is modified with a falling sphere climatology for ALOMAR (Lübken, 1999). In the
forward model we use a linear approximation of the atmosphere with 28 layers, the lowermost located at 22.5 km and the uppermost at 92.5 km. Each layer is 2.5 km thick and no correlation between the layers is introduced mathematically (Rodgers, 1990). However, the altitude resolution of the instrument (see Fig. 1) is larger than the layer thickness in the forward model thus each layer will be affected by layers directly above and below. The OEM is a statistical method of combing two (or more) measurements in an optimal way. As the second set of measurements an estimate of a standard water vapour profile, called the a priori profile, and its associated covariance matrix, is used. This observation focuses on the short term variability and less on the absolute quantities of the water vapour and therefore a static a priori profile was used. It is a piecewise linear profile, increasing from 4 ppmv at 20 km to 7.5 ppmv at 40 km, constant with altitude up to 65 km where it starts to decrease and at 90 km there is no water left (0 ppmv). The same profile has been used for earlier work with WASPAM-data Seele and Hartogh (1999, 2000). An estimate of the quality of the data can be seen from the averaging kernels shown in Fig. 1. Averaging kernels are not expected to change significantly between subsequent measurements as long as the measurement itself has a reasonable signal to noise ratio. The kernels shown can therefore be assumed to be representative for the whole dataset.

Although after six hours of integration time the spectral noise is small enough to retrieve vertical profiles with well defined error bars, the signal-to-noise ratio is still not high enough to detect small tide amplitudes in the water vapour, especially in the upper mesosphere. To overcome this problem we employ a moving time-frame integration scheme. The initial six hour integration time is kept, but every single day in the month is added to the integration. This adds up to a total of \( \sim 180 \) h of integration time for the each monthly time-frame spectra. The timestamp is defined as the time of the day plus/minus half the integration time. After the integration run the time-frame is shifted two hours – if the first data point was between 00:00 and 06:00 UT, the next one will be 02:00 and 08:00 UT. Thus we get 12 spectra for the monthly mean 24 h period. As can be seen in Fig. 1 the averaging kernels from the instrument is almost identical between 50 and 80 km altitude. In other words the amount of information taken from the a priori relative to the observation is very similar.

4 Observations and analysis

Since we only present observations from one location, the source and structure of the oscillations cannot be uniquely deduced. Although we can not conclusively attribute the detected oscillations to the solar tides we have for simplicity decided to denote the 24 h period oscillation as the diurnal component and the 12 h period oscillation as semidiurnal component of the tides. To fit the observed variations a simple wave-like behaviour is assumed,

\[
B + \sum A_i \cdot \cos(\omega_i + \lambda_i)
\]

(1)

where \( B \) denotes the background water vapour mixing ratio, \( A_i \) is the amplitude, \( \omega_i \) the frequency and \( \lambda_i \) the phase offset of the tide. Initially three components where included in the fit (diurnal, semidiurnal and terdiurnal, \( i = 1, 2, 3 \)), however the terdiurnal component was found to be negligible and removed.

In Fig. 2 the WVMR from December 2008 for three different layers can be seen. The altitude difference between each layer (each subplot) is 10 km. The solid line is the measured variability of the WVMR as a function of the local time, and the dotted line is the fit according to Eq. (1). The chosen month for Fig. 2 is arbitrary but shows a general behaviour of the variability of the retrieved volume mixing ratio. Over the course of the year the phase and amplitude is not stable but changes. The variation of the amplitude for the diurnal component can be seen in Fig. 3 and the semidiurnal component in Fig. 4. The detected amplitudes are most probably underestimated as the integration of six hours from every day in the month will smooth out peaks and hereby reduce the amplitude.
Fig. 2. The volume mixing ratio as a function of local time for the eight topmost layers during December 2008. The vertical difference between each layer is 10 km. The solid line is the measured variability and the dotted line is the fit according to Eq. (1.)

Fig. 3. A contour plot showing the variability over the dataset of the diurnal component of the fitted tide. Note the strong increase in the amplitude at each equinox.

Fig. 4. As Fig. 3, but for the semidiurnal component.
Fig. 5. The same state of affairs as Fig. 3, but showing the phase.

The amplitude of the diurnal component shows a semianual variation, with maxima during the equinox and minima during solstice. For the semidiurnal component (Fig. 4) the picture is different with a much more complex behaviour in the measured region. Thus a larger data-set is needed to fully isolate yearly recurring features. In general the amplitude of the diurnal component is stronger than the semidiurnal component. They are however much smaller than the seasonal variability which at 70 km altitude varies with almost a factor of three (Seele and Hartogh, 1999).

The phases (see Figs. 5 and 6) of the observed tidal components changes over the course of the year, which also can be observed in wind fields presented in Portnyagin (2006). Similar to the theoretical results for wind and temperature fields in Forbes (1982a,b) they are relatively stable over the observed altitude with only slight variations for each time step. The semidiurnal component is however more variable than the diurnal counterpart.

5 Discussion

It can be seen from the averaging kernels (Fig. 1) that the atmosphere is fairly well resolved up to 80 km. Information about the water vapour above this level can only be understood as the total column depth which is projected onto the 80 km layer. This could explain the sometimes irregular behaviour at this altitude. It is important to note that the measured quantity is water vapour, an atmospheric tracer. At the altitudes observed water vapour has a lifetime on the order of days and longer, and we can therefore exclude any tidal influence in the water vapour production (or destruction). The observed variability is therefore assumed to be caused by tidal behaviour in the transport mechanisms. Tidal behaviour has been observed in all wind fields (zonal, meridional and vertical) which implies that the resulting tidal transport vector for a tracer would be the sum of the components and as expected the observed behaviour does not follow any of the wind fields.

The seasonal behaviour of the diurnal tide with maxima during equinoxes has been noted by other groups: in global observations of winds from the UARS (Upper Atmosphere Research Satellite) (Hays et al., 1994), as well as ground based radar observations (Vincent et al., 1988) and models (McLandress, 1997). The behaviour is further investigated in two papers by McLandress (2002a,b). In the first paper it is ruled out that the seasonal variation is caused by gravity or planetary wave-tide interactions. Instead, as is shown in the second paper, it is most probably a combination of tidal heating and the zonal mean winds. The model produces qualitatively similar results for both the heating and the mean winds and both effects are therefore used to explain the variation. The effect of the tidal heating can be understood by decomposition of the tidal temperature into its Hough modes. The (1,1) and (1,2) mode are the two dominant vertically propagating modes and due to their different vertical wavelengths constrictive and destructive interference occurs depending on altitude with the resulting annual variation. It should however be noted that the result from these model calculations have not been verified for the high latitudes where the current observations take place.

The retrieval is sensitive to changes in the temperature of the background atmosphere and oscillating temperatures could introduce artifacts in the observations. Currently, the background atmosphere used has a time-step of one day, i.e. each 24 h period is retrieved with the same background. In order to quantify the impact of unaccounted thermal oscillations a sine-like variation of the temperature was modeled and thereafter retrieved with a static background. Oscillations with an amplitude weaker than ±5 K (at 75 km) were found to be negligible. Recent LIDAR results from
ALOMAR (Lübken et al., 2011) indicate thermal tides of a similar magnitude approximately 10 km higher, in the region between 85 and 96 km, which would indicate smaller amplitudes at 75 km. It should also be noted that results from GSWM-09 (Zhang et al., 2010a,b) indicate much weaker thermal tides in this region (±1 K). We therefore assume that the unaccounted thermal tides do not introduce any large artifacts in the data. Oscillations in the background atmosphere do nevertheless introduce a certain small error in the retrieval. Thus a background atmosphere with higher time-resolution would be desirable for a better understanding of the tidal transport mechanisms in the middle atmosphere.

6 Conclusions

A new radiometer has been installed at the ALOMAR observatory. The sensitivity has been greatly improved compared to the older WASPAM instrument, which allows for a higher temporal and vertical resolution. We have looked for tidal signatures in data retrieved since May 2008 by applying a moving time-frame integration scheme. A harmonic approach with a diurnal and semidiurnal component was thereafter used in order to fit the resulting variability in water vapour levels. The observations indicate that the diurnal amplitude component in water vapour has a maximum around the equinoxes which is consistent with model predictions for tidal behaviour in winds. The fit of the semidiurnal component on the other hand is less clear and no obvious pattern could be revealed with the current dataset. In order to better understand its behaviour more data is needed.

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