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*Supplement of*

## **Fine particle pH and gas–particle phase partitioning of inorganic species in Pasadena, California, during the 2010 CalNex campaign**

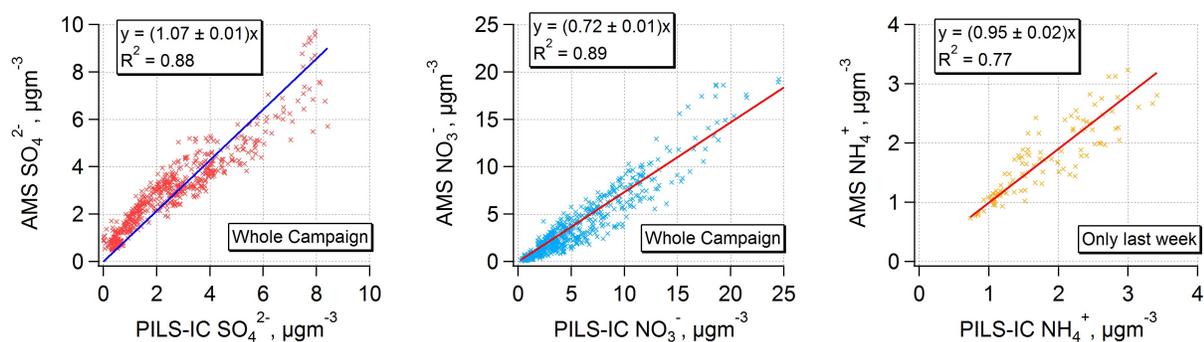
**Hongyu Guo et al.**

*Correspondence to:* Rodney J. Weber ([rweber@eas.gatech.edu](mailto:rweber@eas.gatech.edu)) and Athanasios Nenes ([athanasios.nenes@gatech.edu](mailto:athanasios.nenes@gatech.edu))

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### 1. Comparisons of inorganic species measurements between AMS (PM<sub>1</sub>) and PILS-IC (PM<sub>2.5</sub>)

Consistency ( $R^2 \geq 0.8$ ) between AMS and PILS-IC was observed. AMS measured nominally PM<sub>1</sub>, whereas PILS-IC measured PM<sub>2.5</sub>. These results are similar with other inter-comparisons reported elsewhere (Hayes et al., 2013). A larger difference in slope for nitrate than sulfate is thought to be due to higher nitrate concentrations in the 1 to 2.5  $\mu\text{m}$  size range. PM<sub>1</sub>/PM<sub>2.5</sub> mass ratios, reported in the main text, differ from slopes shown below due to differences in contributions of lower concentrations to these parameters (ratio vs. slope).

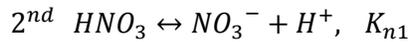
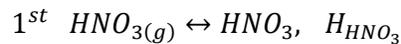


**Fig. S1.** Comparisons of PM<sub>1</sub> AMS sulfate, nitrate, ammonium to PM<sub>2.5</sub> PILS-IC (complete CalNex study except ammonium only for last week). Orthogonal distance regression (ODR) fits with fixed zero intercepts were applied. Fit slope uncertainty is one standard deviation.

## 2. The dependencies of nitrate, ammonium, and chloride on pH, $W_i$ , and T (S curve equation derivations)

### 2.1 $\text{HNO}_3\text{-NO}_3^-$ partitioning

The S curve of  $\varepsilon(\text{NO}_3^-)$  has been discussed explicitly and compared to observations from WINTER aircraft campaign in Guo et al. (2016). Here we show the detailed derivation of equation (3) in that paper. Equilibrium between gaseous  $\text{HNO}_3$  and particle-phase  $\text{NO}_3^-$  involves two processes, first dissolution of  $\text{HNO}_3$  into aqueous phase (assuming particles are liquids) and second dissociation of dissolved  $\text{HNO}_3$  into  $\text{H}^+$  and  $\text{NO}_3^-$ . The two processes are reversible and often reach thermodynamic equilibriums at ambient conditions (RH, T) for fine particles.



for which reaction equilibriums are expressed as follows,

$$H_{\text{HNO}_3} = \gamma_{\text{HNO}_3}[\text{HNO}_3]/p_{\text{HNO}_3} \quad (1)$$

$$K_{n1} = \frac{\gamma_{\text{NO}_3^-}[\text{NO}_3^-]\gamma_{\text{H}^+}[\text{H}^+]}{\gamma_{\text{HNO}_3}[\text{HNO}_3]} \quad (2)$$

where  $H_{\text{HNO}_3}$  is  $\text{HNO}_3$  Henry's law constant,  $K_{n1}$  is  $\text{HNO}_3$  acid dissociation constant,  $\gamma$  represents activity coefficient,  $p_{\text{HNO}_3}$  is partial pressure of  $\text{HNO}_3$  in atmosphere, and  $[x]$  represents aqueous concentrations (mole  $\text{L}^{-1}$ ). From equations (1) and (2) we get the total dissolved  $\text{HNO}_3$  or total particle-phase nitrate ( $\text{NO}_3^{\text{T}}$ ) as

$$[\text{NO}_3^{\text{T}}] = [\text{HNO}_3] + [\text{NO}_3^-] = H_{\text{HNO}_3}p_{\text{HNO}_3} \left( \frac{1}{\gamma_{\text{HNO}_3}} + \frac{K_{n1}}{\gamma_{\text{NO}_3^-}\gamma_{\text{H}^+}[\text{H}^+]} \right) \quad (3)$$

Ideal gas law gives

$$c(\text{HNO}_3) = \frac{p_{\text{HNO}_3}}{RT} \quad (4)$$

where  $c(x)$  represents concentration per volume of air (mole  $\text{m}^{-3}$ ). Therefore, the particle-phase fraction of nitrate is

$$\varepsilon(\text{NO}_3^{\text{T}}) = \frac{c(\text{NO}_3^{\text{T}})}{c(\text{HNO}_3) + c(\text{NO}_3^{\text{T}})} = \frac{[\text{NO}_3^{\text{T}}]W_i}{c(\text{HNO}_3) + [\text{NO}_3^{\text{T}}]W_i} \quad (5)$$

where  $W_i$  is the particle liquid water content associated with inorganic species ( $\mu\text{g m}^{-3}$ ; mass per volume of air) (here the organics associated liquid water is not considered, but it can be included, or measured particle water can be used). Taking equations (3) and (4) into (5), we get  $\varepsilon(\text{NO}_3^{\text{T}})$  as

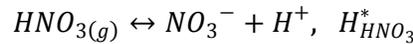
$$\varepsilon(NO_3^T) = \frac{\left(\frac{\gamma_{NO_3^-} \gamma_{H^+}}{\gamma_{HNO_3}} [H^+] + K_{n1}\right) H_{HNO_3} W_i RT}{\gamma_{NO_3^-} \gamma_{H^+} [H^+] + \left(\frac{\gamma_{NO_3^-} \gamma_{H^+}}{\gamma_{HNO_3}} [H^+] + K_{n1}\right) H_{HNO_3} W_i RT} \quad (6)$$

At 298 K,  $K_{n1}=12 \text{ mole L}^{-1}$  (Fountoukis and Nenes, 2007) often  $\gg \frac{\gamma_{NO_3^-} \gamma_{H^+}}{\gamma_{HNO_3}} [H^+]$ , so we assume

$\left(\frac{\gamma_{NO_3^-} \gamma_{H^+}}{\gamma_{HNO_3}} [H^+] + K_{n1}\right) \approx K_{n1}$ . Thus, a simplified equation is

$$\varepsilon(NO_3^T) \cong \frac{K_{n1} H_{HNO_3} W_i RT}{\gamma_{NO_3^-} \gamma_{H^+} [H^+] + K_{n1} H_{HNO_3} W_i RT} \quad (7)$$

$H_{HNO_3} K_{n1}$  is denoted as  $H_{HNO_3}^*$  ( $\text{mole}^2 \text{ kg}^{-2} \text{ atm}^{-1}$ ) hereafter, which is equilibrium constant of the combined dissolution and deprotonation processes as,



$H_{HNO_3}^*$  can be easily calculated by equation (40) in Clegg and Brimblecombe (1990) for T dependence and converted from unit  $\text{atm}^{-1}$  (mole fraction based) to  $\text{mole}^2 \text{ kg}^{-2} \text{ atm}^{-1}$  (molality based) by equation (5) also in that paper. To be consistent with SI units, we have the following equation ready for users' input:

$$\begin{aligned} \varepsilon(NO_3^T) &\cong \frac{H_{HNO_3}^* W_i RT \times 0.987 \times 10^{-14}}{\gamma_{NO_3^-} \gamma_{H^+} [H^+] + H_{HNO_3}^* W_i RT \times 0.987 \times 10^{-14}} \\ &= \frac{H_{HNO_3}^* W_i RT \times 0.987 \times 10^{-14}}{\gamma_{NO_3^-} \gamma_{H^+} 10^{-pH} + H_{HNO_3}^* W_i RT \times 0.987 \times 10^{-14}} \end{aligned} \quad (8)$$

Note that 0.987 comes from the conversion from 1 atm to 1 bar and  $W_i$  unit is  $\mu\text{g m}^{-3}$ . Equation (8) describes the dependence of  $HNO_3$ - $NO_3^-$  partitioning on pH, T, and  $W_i$  (determined by RH and aerosol composition). Based on ideal and non-ideal aqueous particles, several  $\varepsilon(NO_3^-)$  S curves at atmosphere relevant conditions are plotted together with  $\varepsilon(Cl^-)$  and  $\varepsilon(NH_4^+)$  in Fig S3 and S4, respectively.

$\varepsilon(NO_3^T)$  is equivalent to  $\varepsilon(NO_3^-)$  in the main text, since  $NO_3^-$  is practically 100% of  $NO_3^T$  based on  $K_{n1} \gg \frac{\gamma_{NO_3^-} \gamma_{H^+}}{\gamma_{HNO_3}} [H^+]$  (also under atmospheric condition). The fraction of  $NO_3^-$  over  $NO_3^T$  can be given as

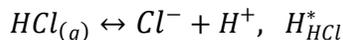
$$\frac{[NO_3^-]}{[NO_3^T]} = \frac{K_{n1}}{K_{n1} + \frac{\gamma_{NO_3^-} \gamma_{H^+}}{\gamma_{HNO_3}} [H^+]} \quad (9)$$

## 2.2 HCl-Cl<sup>-</sup> partitioning

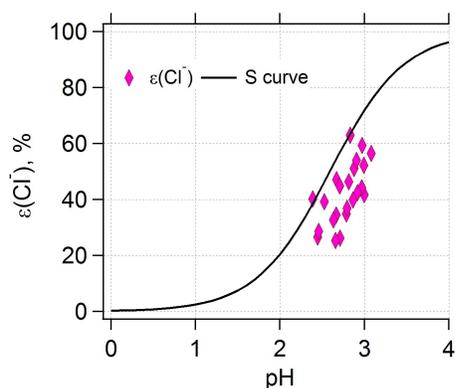
Following the same derivation procedure as  $HNO_3$ - $NO_3^-$  partitioning, we have  $\varepsilon(Cl^-)$  as

$$\begin{aligned}\varepsilon(\text{Cl}^-) &\cong \frac{H_{\text{HCl}}^* W_i RT \times 0.987 \times 10^{-14}}{\gamma_{\text{Cl}^-} \gamma_{\text{H}^+} [\text{H}^+] + H_{\text{HCl}}^* W_i RT \times 0.987 \times 10^{-14}} \\ &= \frac{H_{\text{HCl}}^* W_i RT \times 0.987 \times 10^{-14}}{\gamma_{\text{Cl}^-} \gamma_{\text{H}^+} 10^{-\text{pH}} + H_{\text{HCl}}^* W_i RT \times 0.987 \times 10^{-14}}\end{aligned}\quad (10)$$

where  $H_{\text{HCl}}^*$  (mole<sup>2</sup> kg<sup>-2</sup> atm<sup>-1</sup>) is the equilibrium constant and is equal to the “conventional” Henry’s law constant multiplied by the acid dissociation constant of hydrochloric acid.  $H_{\text{HCl}}^*$  can be calculated by equation (22) in Carslaw et al. (1995) to account for T’s variation.



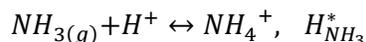
A comparison of  $\varepsilon(\text{Cl}^-)$  S curve with a subset of CalNex data is shown in Fig. S2. The selected CalNex data are all in a small range of T 15.5 to 19.5 °C (around campaign average T) and  $W_i$  10 to 20  $\mu\text{g m}^{-3}$ , while the S curve is calculated based on the average condition of these data as T = 17.5 °C,  $W_i = 15 \mu\text{g m}^{-3}$ ,  $\gamma_{\text{H}^+ - \text{Cl}^-} = 0.81$  (The binary activity coefficient,  $\gamma_{\text{H}^+ - \text{Cl}^-} = \sqrt{\gamma_{\text{H}^+} \gamma_{\text{Cl}^-}}$ ). The distribution of the  $\varepsilon(\text{Cl}^-)$  points close to S curve validates the PM<sub>2.5</sub> pH prediction and demonstrates the usage of S curve.



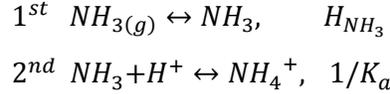
**Fig. S2.** The comparison of S curve and measured  $\varepsilon(\text{Cl}^-)$  with predicted particle pH by ISORROPIA-II.  $\text{Cl}^-$  is from PM<sub>2.5</sub> PILS-IC measurements.

### 2.3 $\text{NH}_3\text{-NH}_4^+$ partitioning

The derivation of  $\text{NH}_3\text{-NH}_4^+$  partitioning is a bit different from the above two acidic gases. Equilibrium between gaseous  $\text{NH}_3$  and  $\text{NH}_4^+$  can be described simply as



( $H_{\text{NH}_3}^*$  is equivalent to the “conventional” Henry’s law constant of  $\text{NH}_3$  divided by the acid dissociation constant of  $\text{NH}_4^+$ ) or described by the follow two reversible reactions assuming water activity as unity.



for which reaction equilibriums are described as

$$H_{\text{NH}_3} = \gamma_{\text{NH}_3} [\text{NH}_3] / p_{\text{NH}_3} \quad (11)$$

$$1/K_a = \frac{\gamma_{\text{NH}_4^+} [\text{NH}_4^+]}{\gamma_{\text{NH}_3} [\text{NH}_3] \gamma_{\text{H}^+} [\text{H}^+]} \quad (12)$$

where  $H_{\text{NH}_3}$  is  $\text{NH}_3$  Henry's law constant,  $K_a$  is  $\text{NH}_4^+$  acid dissociation constant,  $\gamma$  represents activity coefficient,  $p_{\text{NH}_3}$  is partial pressure of  $\text{NH}_3$  in atmosphere, and  $[x]$  represents aqueous concentrations (mole  $\text{L}^{-1}$ ). Please note that the 2<sup>nd</sup> reaction is usually written in another form (Fountoukis and Nenes, 2007) as



where  $K_w$  is water dissociation constant. Equations (10) and (11) give the total dissolved  $\text{NH}_3$  or total particle-phase ammonium ( $\text{NH}_4^{\text{T}}$ ) as

$$[\text{NH}_4^{\text{T}}] = [\text{NH}_3] + [\text{NH}_4^+] = H_{\text{NH}_3} p_{\text{NH}_3} \left( \frac{1}{\gamma_{\text{NH}_3}} + \frac{\gamma_{\text{H}^+} [\text{H}^+]}{\gamma_{\text{NH}_4^+} K_a} \right) \quad (13)$$

Combining with ideal gas law, that is

$$c(\text{NH}_3) = \frac{p_{\text{NH}_3}}{RT} \quad (14)$$

where  $c(x)$  represents concentration per volume of air (mole  $\text{m}^{-3}$ ). We have the particle-phase fraction of ammonium as

$$\varepsilon(\text{NH}_4^{\text{T}}) = \frac{c(\text{NH}_4^{\text{T}})}{c(\text{NH}_3) + c(\text{NH}_4^{\text{T}})} = \frac{[\text{NH}_4^{\text{T}}] W_i}{c(\text{NH}_3) + [\text{NH}_4^{\text{T}}] W_i} \quad (15)$$

With equations (12) and (13), the above equation is transformed into

$$\varepsilon(\text{NH}_4^{\text{T}}) = \frac{\left( \frac{\gamma_{\text{H}^+} [\text{H}^+]}{\gamma_{\text{NH}_4^+}} + \frac{K_a}{\gamma_{\text{NH}_3}} \right) \frac{H_{\text{NH}_3}}{K_a} W_i RT}{1 + \left( \frac{\gamma_{\text{H}^+} [\text{H}^+]}{\gamma_{\text{NH}_4^+}} + \frac{K_a}{\gamma_{\text{NH}_3}} \right) \frac{H_{\text{NH}_3}}{K_a} W_i RT} \quad (16)$$

At 298 K,  $K_a = 5.69 \times 10^{-10}$  mole  $\text{L}^{-1}$  (Clegg et al., 1998) results in  $\frac{K_a}{\gamma_{\text{NH}_3}} \ll \frac{\gamma_{\text{H}^+} [\text{H}^+]}{\gamma_{\text{NH}_4^+}}$  as long as the solution is

not too basic. Neglecting  $\frac{K_a}{\gamma_{\text{NH}_3}}$  part and taking  $\frac{H_{\text{NH}_3}}{K_a} = H_{\text{NH}_3}^*$ , we have

$$\varepsilon(\text{NH}_4^{\text{T}}) \cong \frac{\frac{\gamma_{\text{H}^+} [\text{H}^+]}{\gamma_{\text{NH}_4^+}} H_{\text{NH}_3}^* W_i RT}{1 + \frac{\gamma_{\text{H}^+} [\text{H}^+]}{\gamma_{\text{NH}_4^+}} H_{\text{NH}_3}^* W_i RT} \quad (17)$$

To be consistent with SI units, the equation (16) is then presented as:

$$\begin{aligned} \varepsilon(\text{NH}_4^{\text{T}}) &\cong \frac{\frac{\gamma_{\text{H}^+}[\text{H}^+]}{\gamma_{\text{NH}_4^+}} H_{\text{NH}_3}^* W_i RT \times 0.987 \times 10^{-14}}{1 + \frac{\gamma_{\text{H}^+}[\text{H}^+]}{\gamma_{\text{NH}_4^+}} H_{\text{NH}_3}^* W_i RT \times 0.987 \times 10^{-14}}} \\ &= \frac{\frac{\gamma_{\text{H}^+} 10^{-\text{pH}}}{\gamma_{\text{NH}_4^+}} H_{\text{NH}_3}^* W_i RT \times 0.987 \times 10^{-14}}{1 + \frac{\gamma_{\text{H}^+} 10^{-\text{pH}}}{\gamma_{\text{NH}_4^+}} H_{\text{NH}_3}^* W_i RT \times 0.987 \times 10^{-14}} \end{aligned} \quad (18)$$

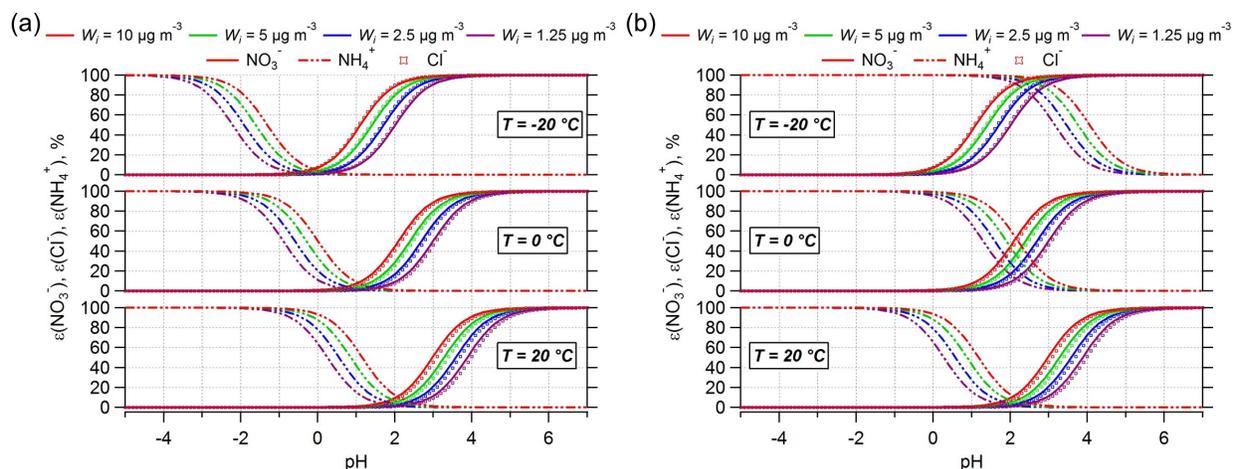
where the 0.987 comes from the conversion from 1 atm to 1 bar and  $W_i$  unit is  $\mu\text{g m}^{-3}$ .  $H_{\text{NH}_3}^*$  ( $\text{atm}^{-1}$ ) can be calculated from equation (12) in Clegg et al. (1998) following a typo correction to the equation. The corrected equation is:

$$\ln(H_{\text{NH}_3}^*) = 25.393 - 10373.6(1/T_r - 1/T) + 4.131(T_r/T - (1 + \ln(T_r/T))) \quad (19)$$

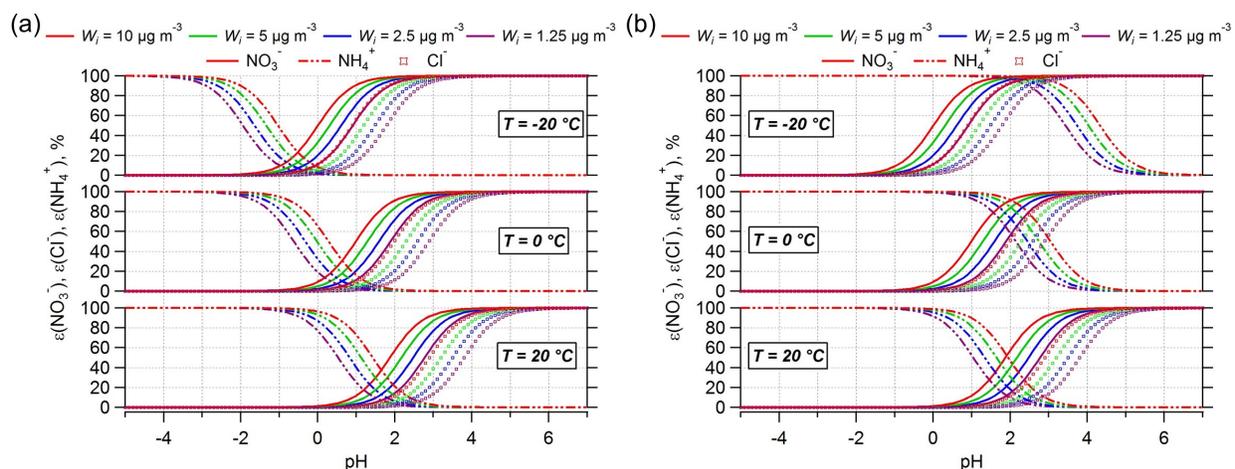
where  $T_r$  is the reference temperature of 298.15 K. Note that, the mole fraction based  $H_{\text{NH}_3}^*$  has the same numerical value as its molality based form. After correction, a larger  $H_{\text{NH}_3}^*$  is associated with a lower temperature, consistent with a general expectation that  $\text{NH}_3$  condenses onto particles if temperature decreases.  $\varepsilon(\text{NH}_4^{\text{T}})$  is equivalent to  $\varepsilon(\text{NH}_4^+)$  presented in the main text, since  $\text{NH}_4^+$  is the dominant form of dissolved  $\text{NH}_3$  based on  $\frac{K_a}{\gamma_{\text{NH}_3}} \ll \frac{\gamma_{\text{H}^+}[\text{H}^+]}{\gamma_{\text{NH}_4^+}}$  and under atmospheric conditions.

**Summary:** with the equations of  $\varepsilon(\text{NO}_3^-)$ ,  $\varepsilon(\text{Cl}^-)$ , and  $\varepsilon(\text{NH}_4^+)$ , S-shaped curves of these three paired gas to particle partitioning can be easily calculated with pH, T,  $W_i$ , and activity coefficients. We simulate two set of results, Fig S3 assuming activity coefficients to be one (ideal solution) and Fig S4 with practical activity coefficients from CalNex.

## Fine particle pH and gas-particle phase partitioning of inorganics during the 2010 CalNex campaign



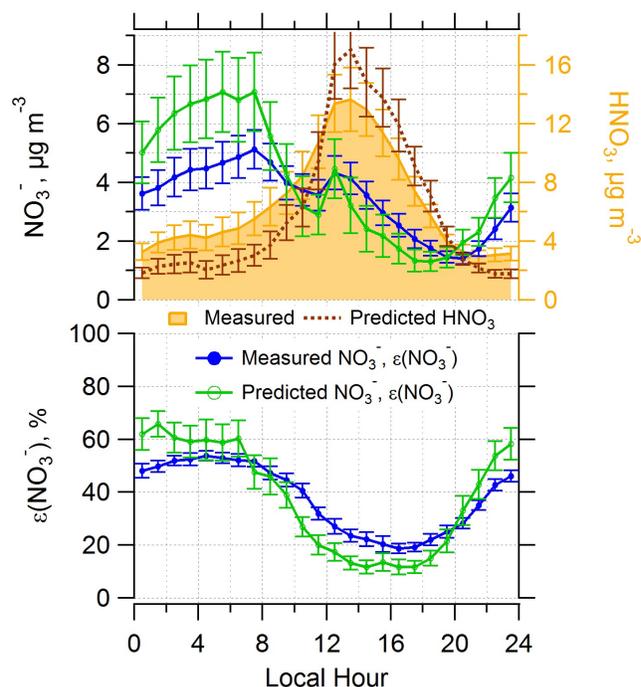
**Fig. S3.** Simulated  $\epsilon(\text{NO}_3^-)$ ,  $\epsilon(\text{NH}_4^+)$ ,  $\epsilon(\text{Cl}^-)$  at  $-20\text{ }^{\circ}\text{C}$ ,  $0\text{ }^{\circ}\text{C}$ ,  $20\text{ }^{\circ}\text{C}$  and various particle liquid water levels ( $1.25$ ,  $2.5$ ,  $5$ ,  $10\text{ }\mu\text{g m}^{-3}$ ) assuming ideal solutions. Note that Fig. S3a shows  $\epsilon(\text{NH}_4^+)$  calculated using the equation (12) in Clegg et al. (1998), and Fig. S3b shows  $\epsilon(\text{NH}_4^+)$  using the corrected equation, given above as equation (19). The correct equation for  $H_{\text{NH}_3}^*$  shifts  $\epsilon(\text{NH}_4^+)$  to higher pH at lower  $T$ , as expected.



**Fig. S4.** Simulated  $\epsilon(\text{NO}_3^-)$ ,  $\epsilon(\text{NH}_4^+)$ ,  $\epsilon(\text{Cl}^-)$  at  $-20\text{ }^{\circ}\text{C}$ ,  $0\text{ }^{\circ}\text{C}$ ,  $20\text{ }^{\circ}\text{C}$  and various particle liquid water levels ( $1.25$ ,  $2.5$ ,  $5$ ,  $10\text{ }\mu\text{g m}^{-3}$ ) with activity coefficients obtained from CalNex campaign.  $\gamma_{\text{H}^+ - \text{NO}_3^-} = 0.28$ ,  $\gamma_{\text{H}^+ - \text{Cl}^-} = 0.81$ , and  $\gamma_{\text{H}^+} / \gamma_{\text{NH}_4^+} = 1.90$ . Similar to Fig. S3, Fig. S4a shows  $\epsilon(\text{NH}_4^+)$  calculated using the equation (12) in Clegg et al. (1998), and Fig. S4b shows  $\epsilon(\text{NH}_4^+)$  using the corrected equation, given above as equation (19).

### 3. Investigation of the cause for bias in $\epsilon(\text{NO}_3^-)$ : sample line heating?

As Fig. S5 shows,  $\text{NO}_3^-$  and  $\epsilon(\text{NO}_3^-)$  are both over-predicted during the nighttime and under-predicted during the daytime. The deviations from measurements are anti-correlated with nitric acid. The deviation between predicted and measured  $\text{HNO}_3$  also has a diurnal pattern, reverse to that of  $\text{NO}_3^-$ .



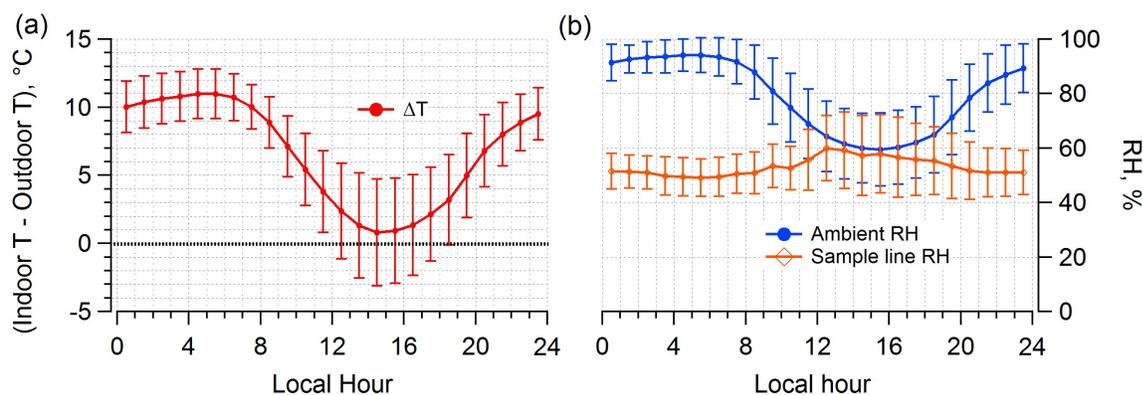
**Fig. S5.** Diurnal profiles of measured and predicted  $\text{HNO}_3$ ,  $\text{NO}_3^-$ , and  $\epsilon(\text{NO}_3^-)$ . Data shown above are for the complete CalNex study and particle-phase data is AMS  $\text{PM}_{10}$ . Mean hourly averages are shown and standard errors are plotted as error bars.

**Table S1.** Summary of temperature differences in sample lines and ambient and sample line residence time for the AMS and CIMS. AMS indoor T was 25°C. CIMS inlet was heated to 75°C.

Instrument	Inlet residence time, sec	Time of the day	Temperature differences, °C
AMS	2.1	Day	~0
		Night	~ +10
CIMS	0.32	Day	~ +50
		Night	~ +60

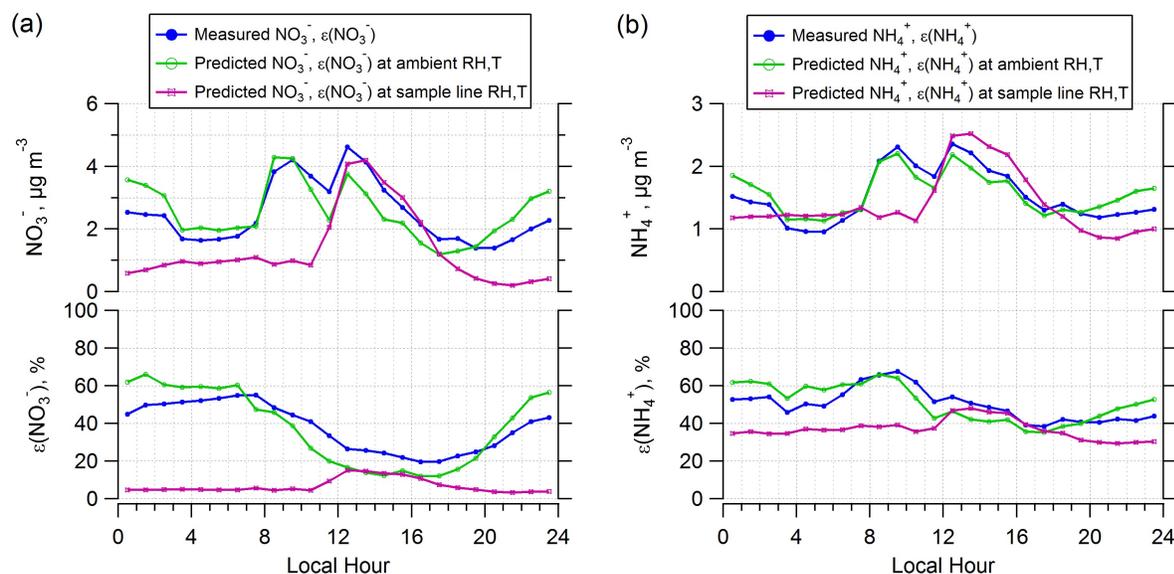
For the AMS sample line located indoors, particle heating was most likely to occur at night (indoor T > ambient T), which may cause semi-volatile NO<sub>3</sub><sup>-</sup> loss. There were no temperature differences during the day (Table S1). To examine the possible sample line heating/cooling effect, we first determined sample line RH (Equation 20) by conservation of water vapor under isobaric condition and following saturated water vapor pressure equation  $e_s = 6.11 \times 10^{\left(\frac{7.5T}{237.5+T}\right)}$  (T unit as °C) (Alduchov and Eskridge, 1996). The inferred sample line RH is plotted with measured ambient RH in Fig. S6b. Sample line RH was lower (~50%) than ambient (~90%) at midnight and close to ambient (~60%) in the afternoon since temperatures were similar.

$$RH_2 = RH_1 10^{\left[\left(\frac{7.5T_1}{237.5+T_1}\right) - \left(\frac{7.5T_2}{237.5+T_2}\right)\right]} \quad (20)$$



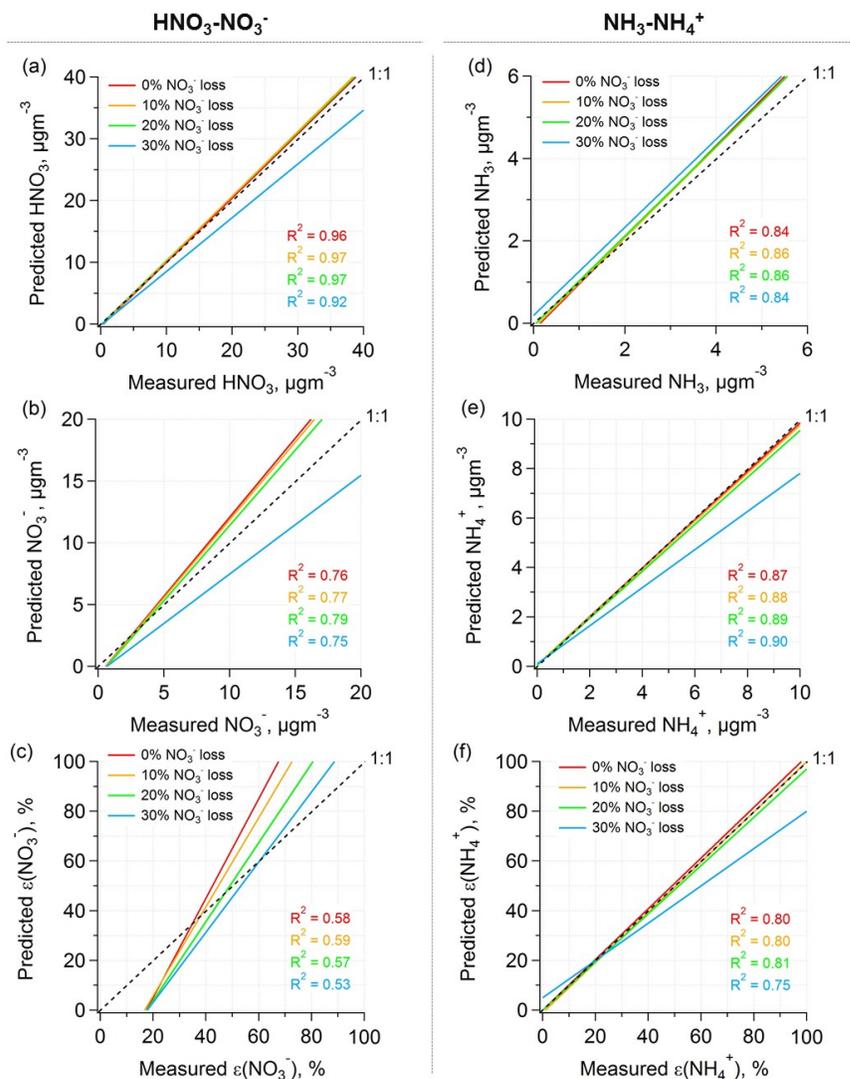
**Fig. S6.** Diurnal profiles of (a) temperature difference between AMS indoor and outdoor and (b) corresponding ambient and RH predicted in the sample line due to the T difference. Mean hourly averages and standard deviations are shown.

ISORROPIA-II was run with aerosol and gas-phase species at the AMS sample line T and RH and compared to predictions from ambient T and RH and measurements. Fig. S7 is discussed in the main text section 4.1.



**Fig. S7.** Diurnal profiles of measured and predicted (a)  $\text{NO}_3^-$ ,  $\epsilon(\text{NO}_3^-)$  and (b)  $\text{NH}_4^+$ ,  $\epsilon(\text{NH}_4^+)$ . Predictions are based on ambient or sample line RH and T for AMS inlet. Data shown above are for the complete CalNex study in the 20-95% RH range and particle-phase data is AMS  $\text{PM}_{10}$ . Mean hourly averages are shown. ISORROPIA run with ambient data show that the predicted partitioning between the particle and gas phase is in better agreement with observations than runs using sample line T and RH. Note that in both runs, only T and RH differ since total nitrate and ammonium input are the same.

CIMS inlet heating is similar for day ( $\sim 50^\circ\text{C}$ ) and night ( $\sim 60^\circ\text{C}$ ). Potential bias in the  $\text{HNO}_3$  or  $\text{HCl}$  then mainly depends on the mass loadings of  $\text{NO}_3^-$  or  $\text{Cl}^-$ . Here we focus only on the possible bias due to over-measurement of  $\text{HNO}_3$ . ISORROPIA-II was run at ambient RH and T with a “corrected”  $\text{HNO}_3$  at three assumed lower levels of  $\text{HNO}_3$  to compensate for an assumed positive nitrate artifact of 10%, 20%, 30% (i.e., assuming 10, 20 or 30% of the nitrate measured by the AMS or PILS was evaporated in the CIMS inlet leading to an over-measurement of  $\text{HNO}_3$ . 10% to 30% particle  $\text{NO}_3^-$  was subtracted from the measured CIMS  $\text{HNO}_3$ ). Only  $\text{HNO}_3$  is modified, all other inputs are kept the same. Results are shown in Fig. S8. Evaporation of 30% of the measured nitrate is expected to be an extreme upper limit. For instance, 66% of  $\text{PM}_{10}$  nitrate evaporated at a temperature of  $75^\circ\text{C}$  in a thermal denuder upstream of the AMS at the CalNex site, consistent with previous results at other urban sites in the LA area and elsewhere (Huffman et al., 2009). The residence time on the thermal denuder was  $\sim 12$  sec, while that of the CIMS inlet was  $\sim 0.32$  sec, so the extent of evaporation in the CIMS inlet assumed to be substantially lower than that in the thermal denuder.



**Fig. S8.** Comparisons of predicted and measured  $\text{HNO}_3$ ,  $\text{NO}_3^-$ , and  $\epsilon(\text{NO}_3^-)$  (a, b, c) and  $\text{NH}_3$ ,  $\text{NH}_4^+$ , and  $\epsilon(\text{NH}_4^+)$  (d, e, f) for data from the complete CalNex study based on “corrected”  $\text{HNO}_3$  data due to assumed  $\text{PM}_{10}$  nitrate evaporation in the heated CIMS inlet. The other inputs are kept the same. Only the orthogonal regression fits are shown. “0%  $\text{NO}_3^-$  loss” condition is the same as Figure 2 in the main text.

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